

Hard Exclusive Reactions

Caroline Riedl



γ

ω

ϕ

π

Υ

ρ

J/ψ

η

ξ

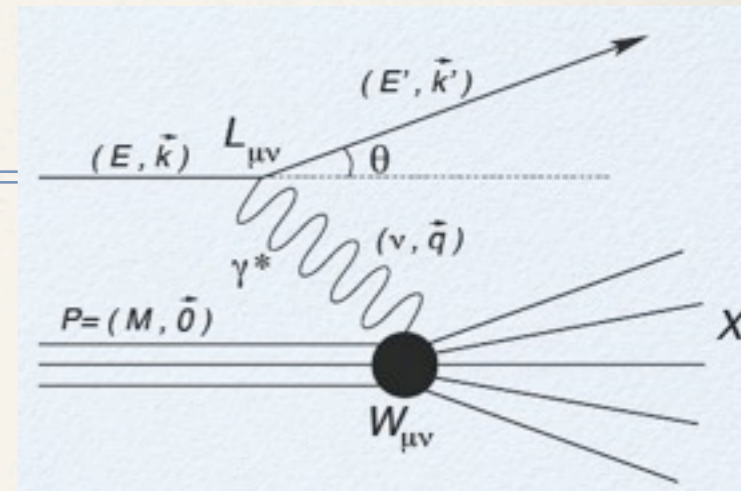
Milos, Greece, 30.9.2009

EINN09: Electromagnetic Interactions with Nucleons and Nuclei

Outline: Hard Exclusive Reactions

Hard scale:
 Q^2 , M , and/or t

Complete spectrum
of X known



● Photoproduction ($Q^2 < 1 \text{ GeV}^2$)

● Cross sections (Q^2, W, t)

testing ground for QCD

● Electroproduction ($Q^2 > 1 \text{ GeV}^2$)

● t slopes

● Azimuthal asymmetries

● Spin density matrix elements

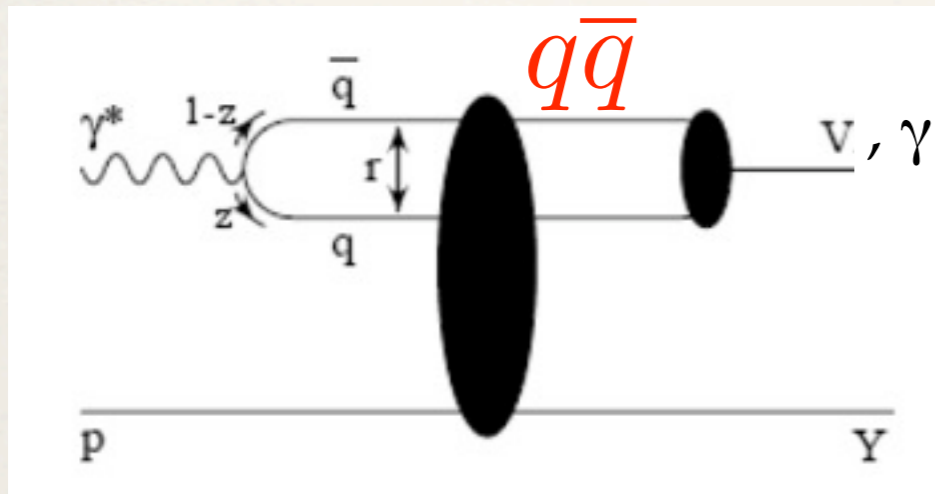
GPDs

total angular
momentum of quarks

● Future

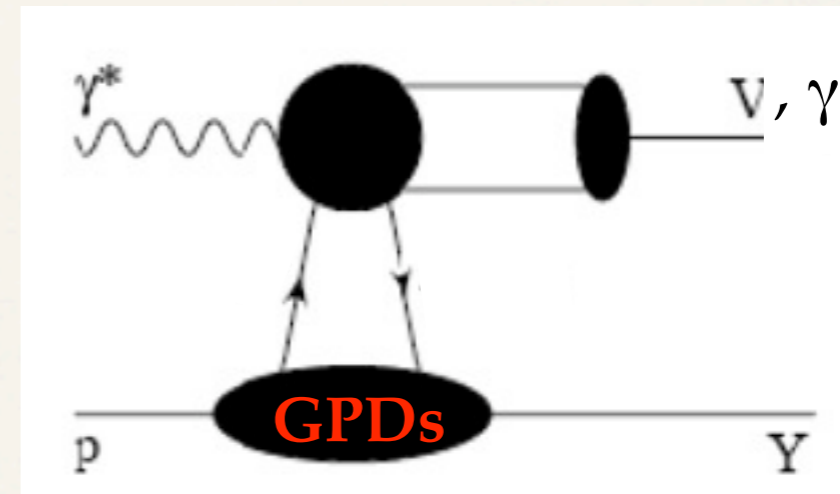
$ep \rightarrow epV$ or $ep \rightarrow ep\gamma$ (DVCS)

High energy factorization



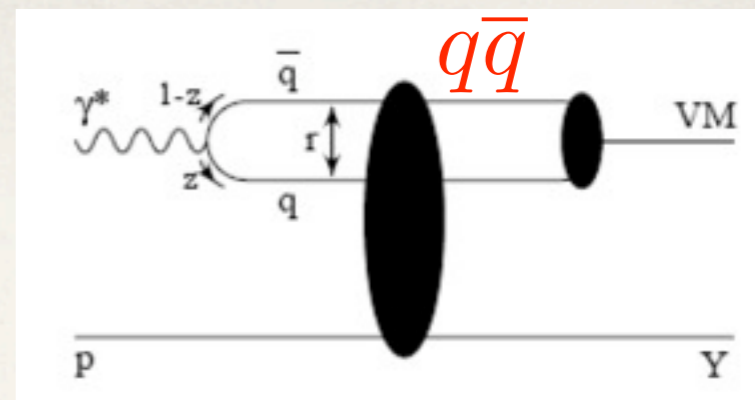
- Universal dipole interactions
- Low $x \leftrightarrow$ large W
- Scale: $Q^2 + m_V^2$

Collinear factorization



- Parameterization of non-perturbative nucleonic structure
- Information on parton-parton correlations
- VM: proven only for $\sigma_{\text{Longitudinal}}$

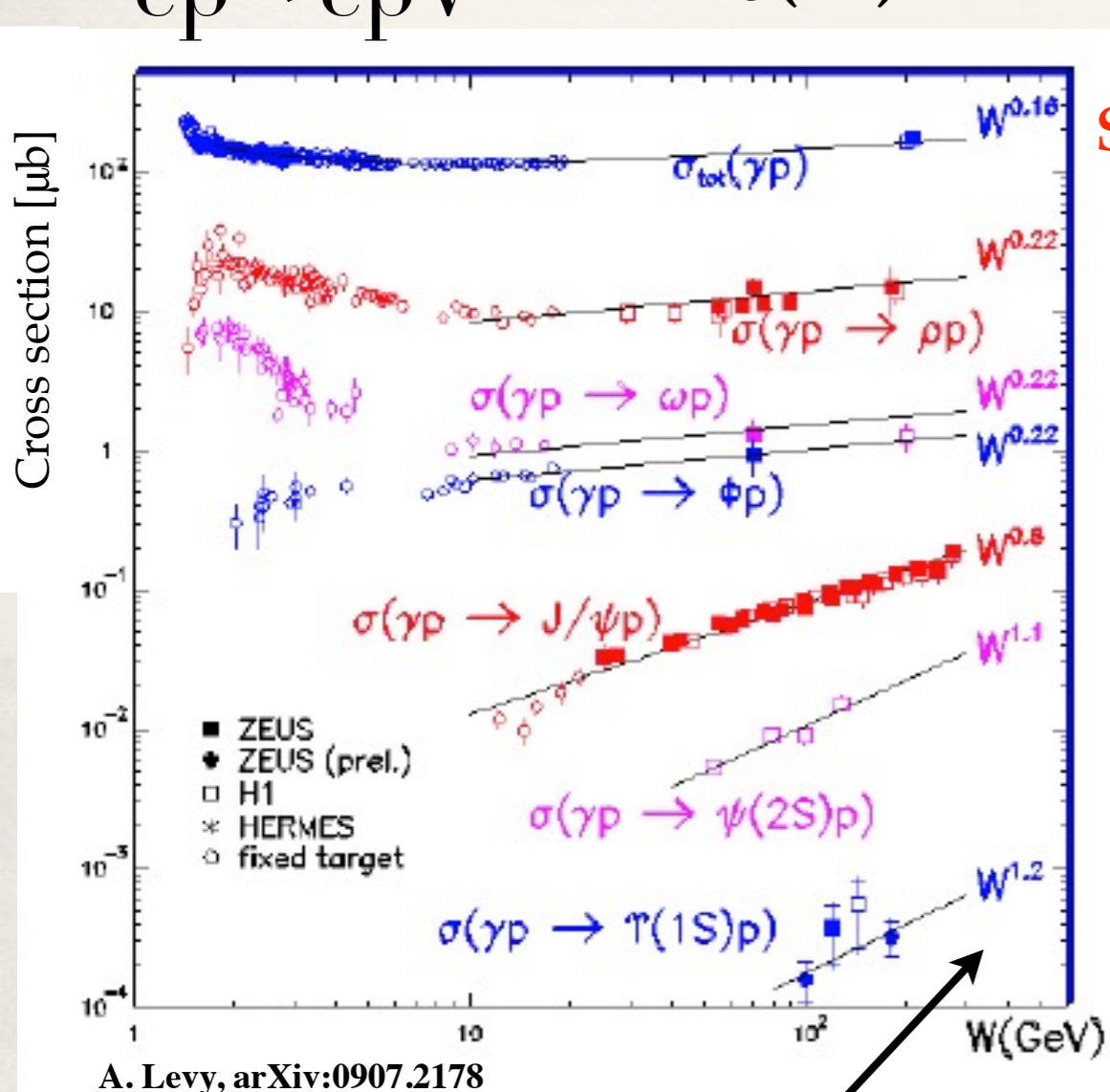
Kinematic Landscape



Photoproduction

$$ep \rightarrow epV$$

$$\sigma(W) \propto W^\delta$$



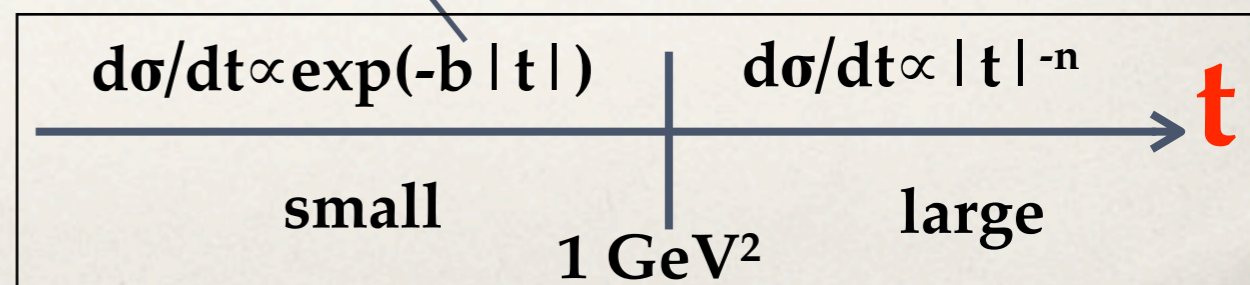
$q\bar{q}$
size

HERA sees interplay between soft and hard

$\delta=0.2$
 Q, ω, ϕ $b=10 \text{ GeV}^{-2}$ **SOFT** Regge, soft pomeron exchange

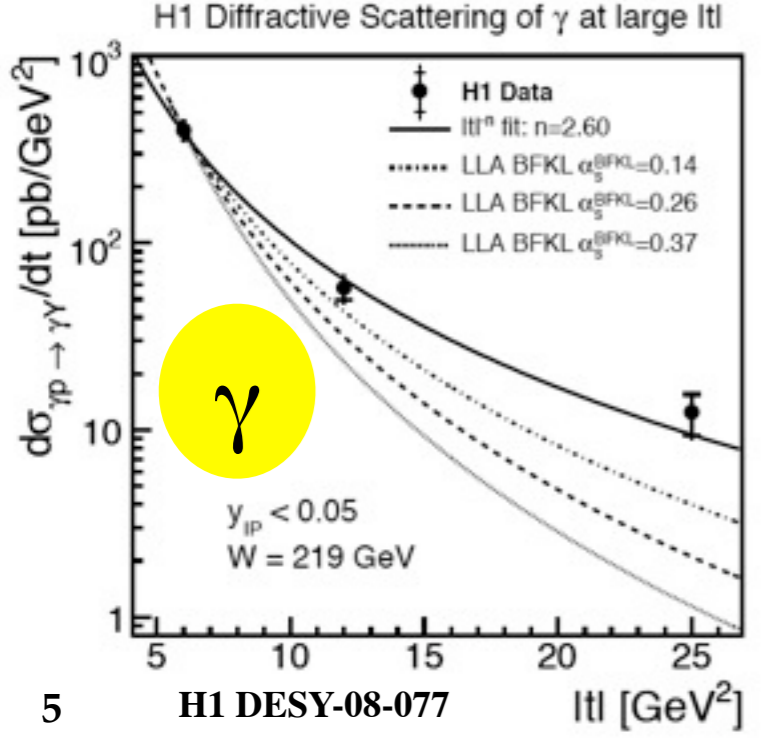
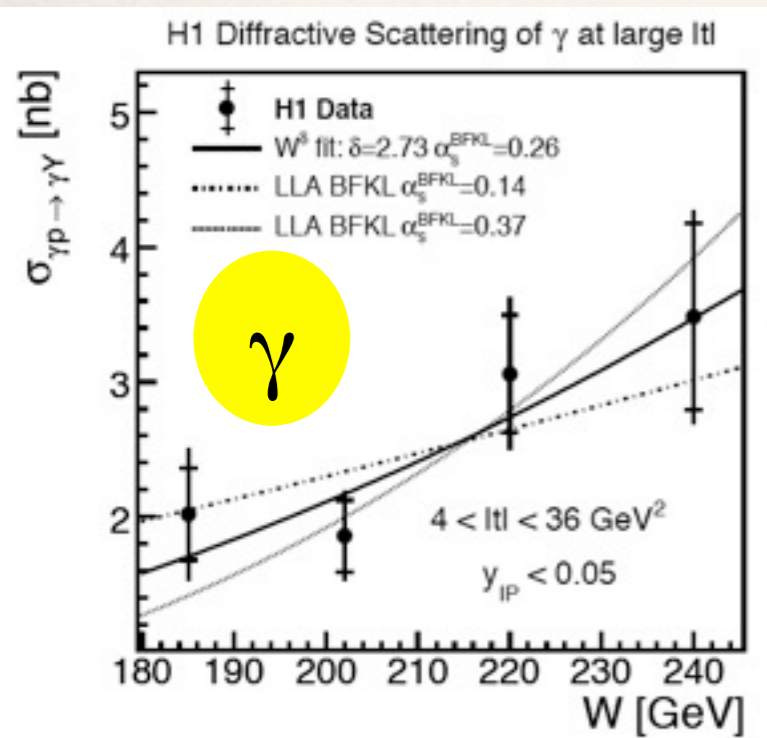
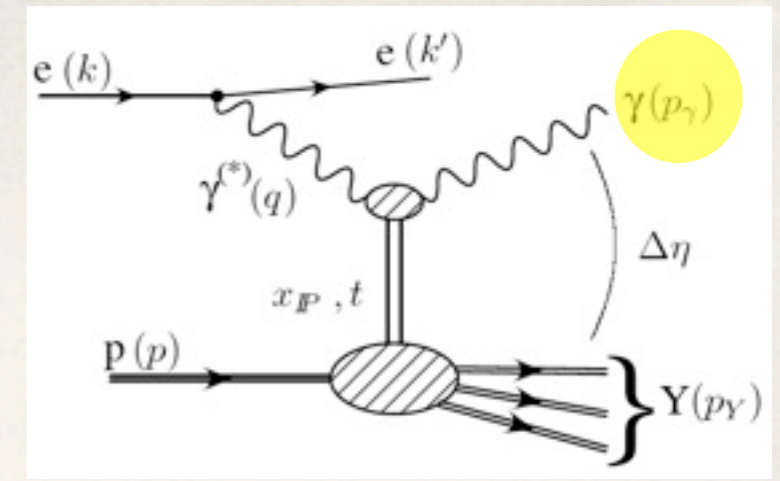
$\delta \geq 0.8$
 $J/\psi, \gamma$ $b=4-5 \text{ GeV}^{-2}$ **HARD** pQCD, 2-gluon exchange

$M_V,$
 $\langle W \rangle$



$$d\sigma/dt \propto [xg(x, Q^2)]^2$$

$\gamma^{(*)} p \rightarrow \gamma Y$ Photoproduction



- Large $|t|$, proton-dissociative

- $\sigma(W) \propto W^\delta$

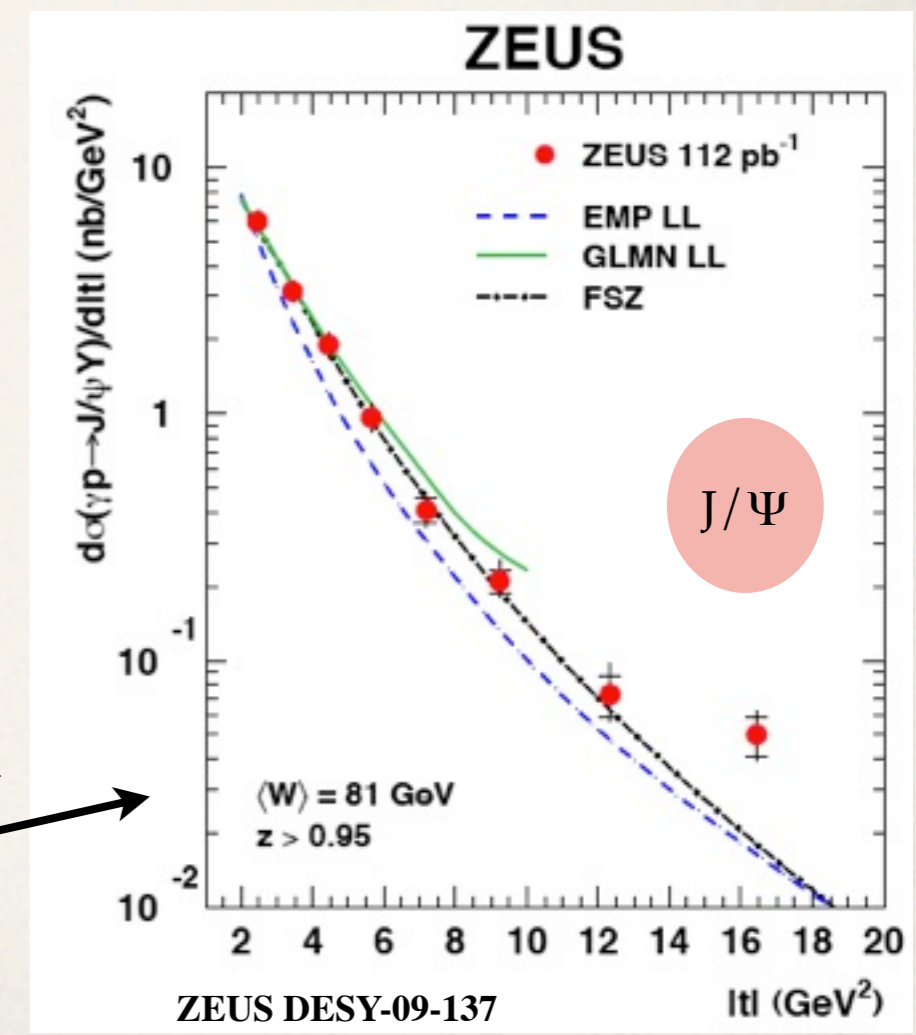
- pQCD: good description

- Compatible with J/Ψ

- $d\sigma/dt \propto |t|^{-n}$

- pQCD: too soft prediction

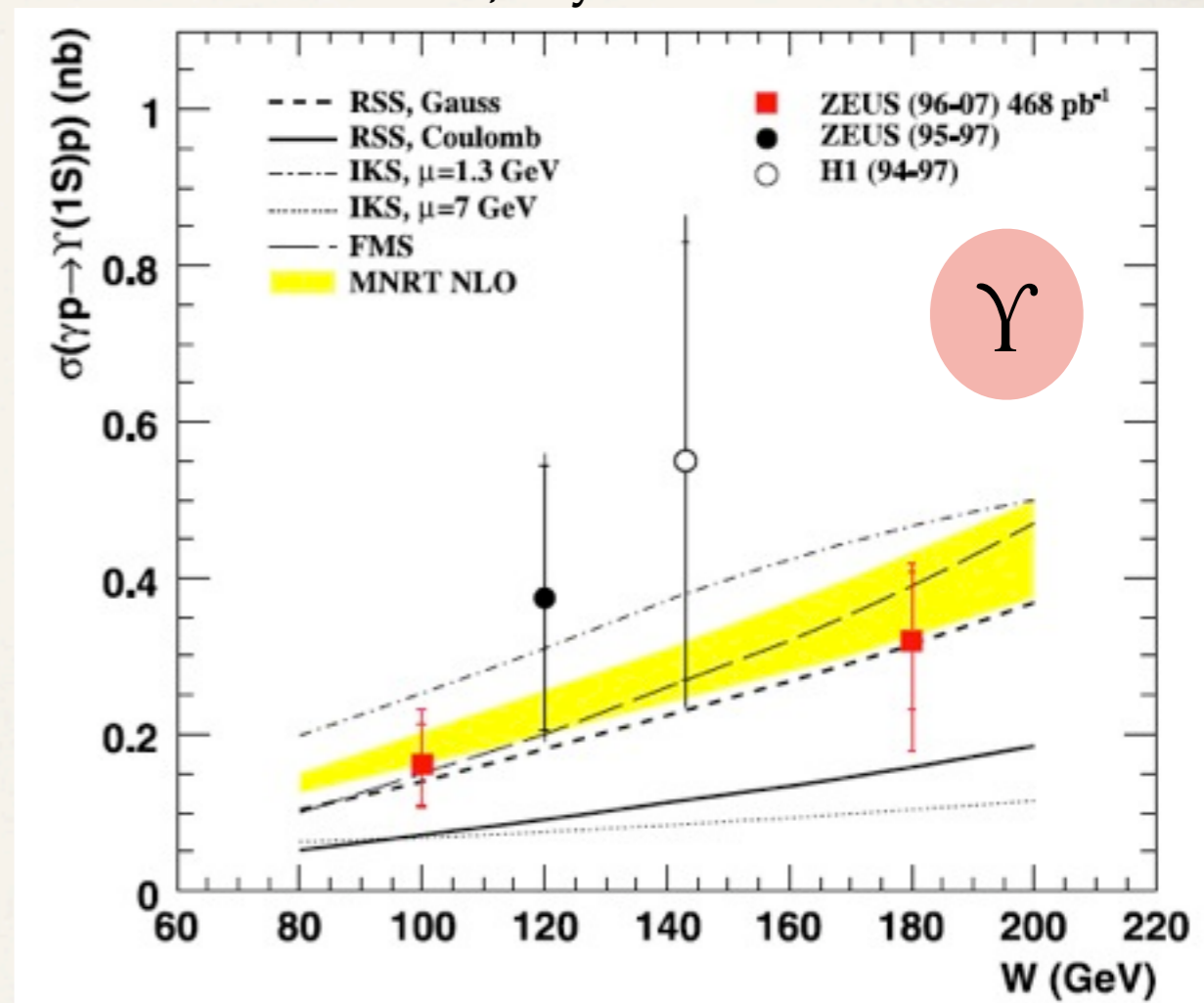
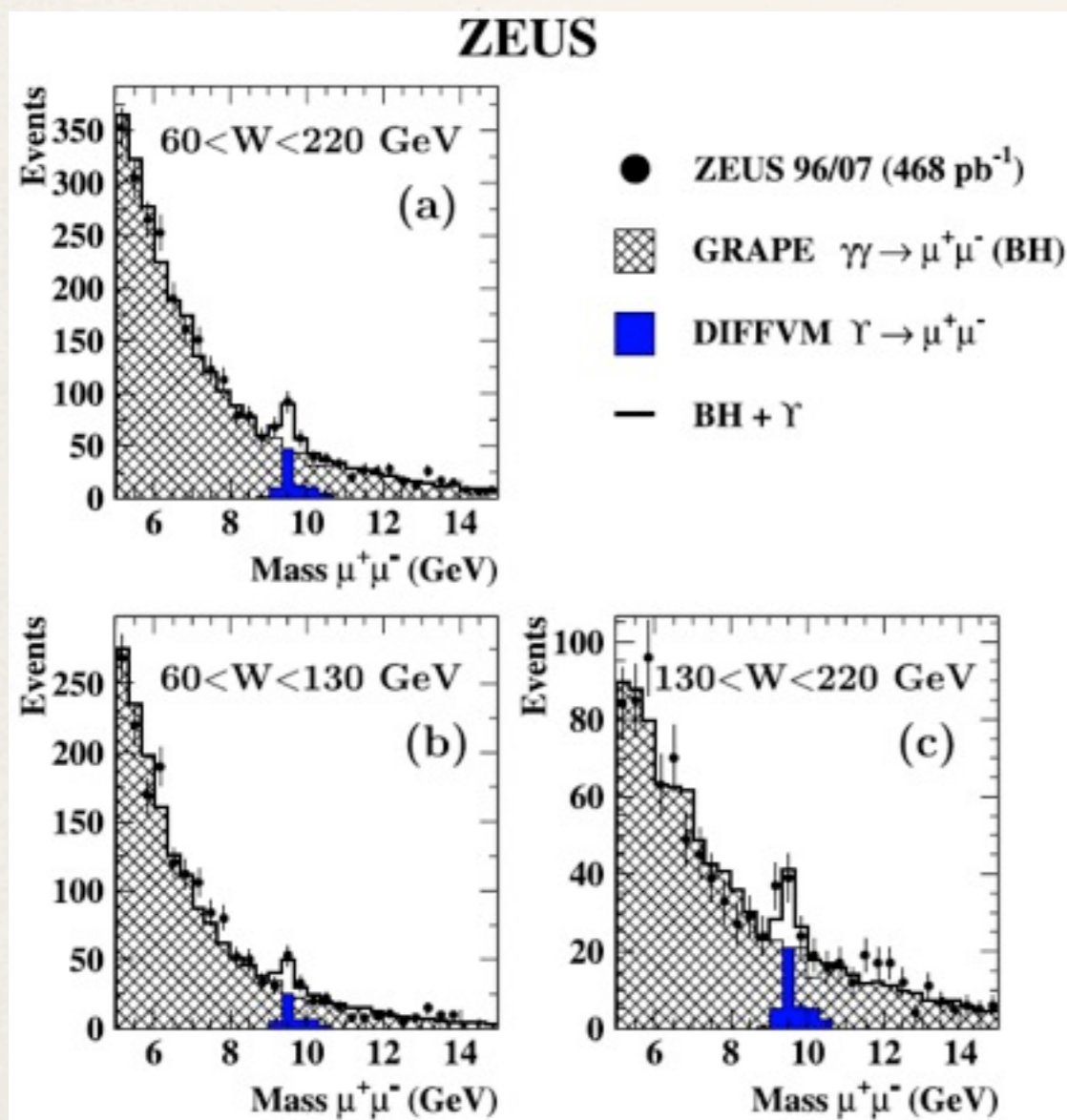
- Harder than J/Ψ



$ep \rightarrow ep \Upsilon(1S)$ Photoproduction

ZEUS 1996-2007, Phys. Lett. B 680 (2009) 4-12

Di-muon events, invariant mass:



$\sigma \propto W^\delta, \delta=1.2 \pm 0.8$

✓ pQCD calculations

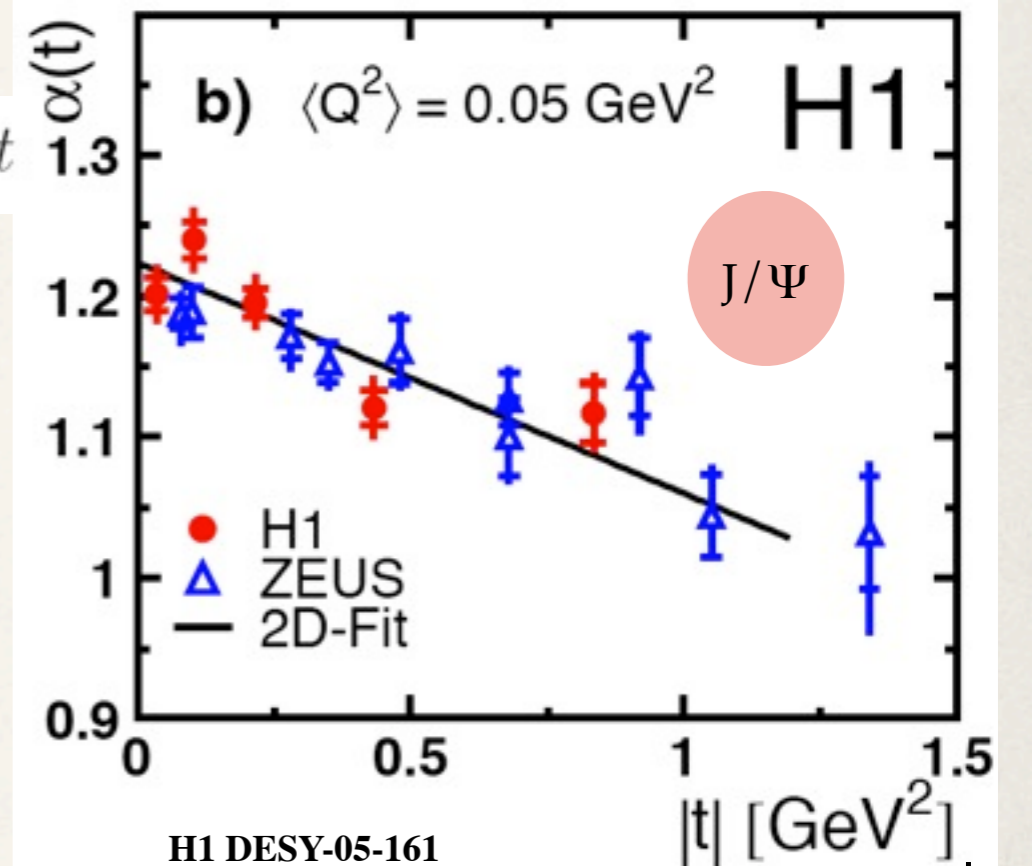
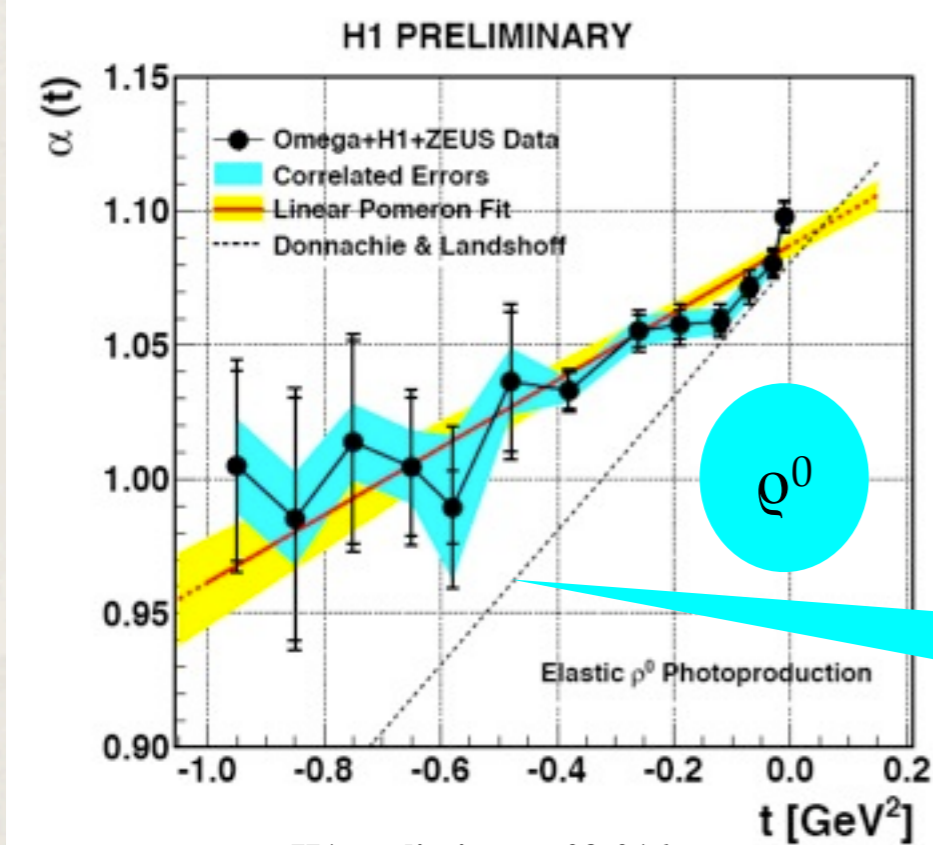
$ep \rightarrow epV$ Pomeron trajectories

$$d\sigma/dt = F(t)W^{4(\alpha_P(t)-1)}$$

$$\alpha_P(t) = \alpha_P(0) + \alpha'_P \cdot t$$

photoproduction light VM

photoproduction heavy VM



Canonical soft pomeron:
 $\alpha'_P = 0.25 \text{ GeV}^{-2}$

$$\alpha'_P = (0.126 \pm 0.013 \pm 0.012) \text{ GeV}^{-2}$$

smaller than canonical soft pomeron

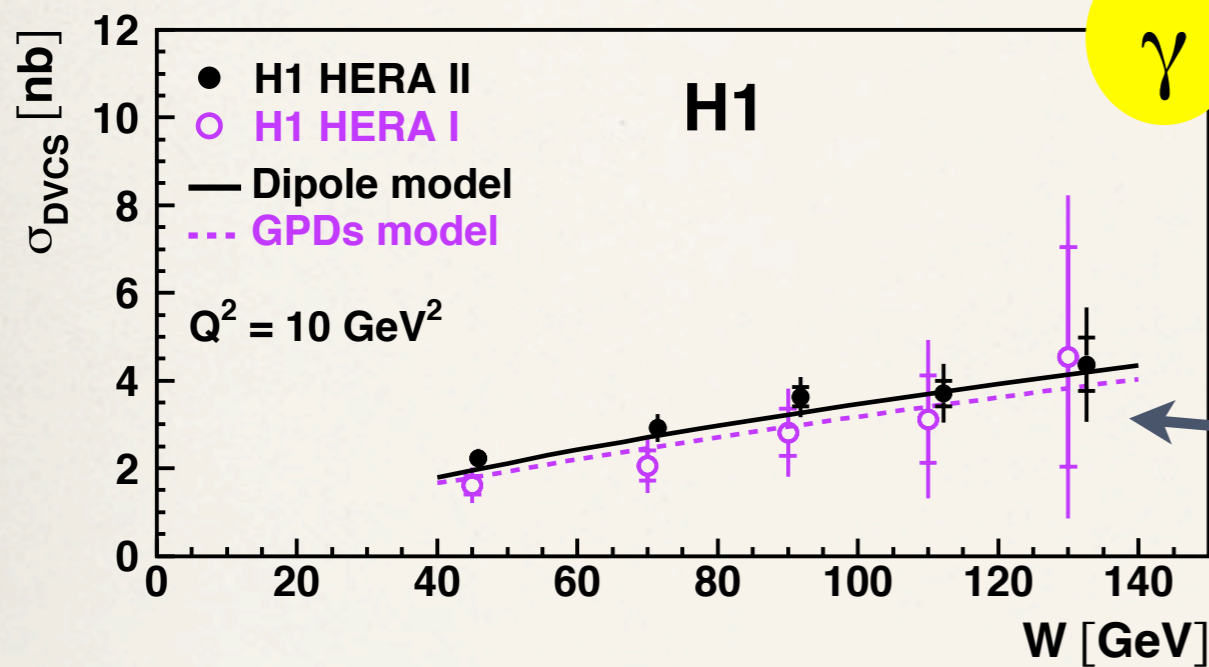
$$\alpha'_P = (0.164 \pm 0.028 \pm 0.030) \text{ GeV}^{-2}$$

hard pomeron known to be smaller than canonical soft pomeron

$ep \rightarrow ep\gamma$ cross-section

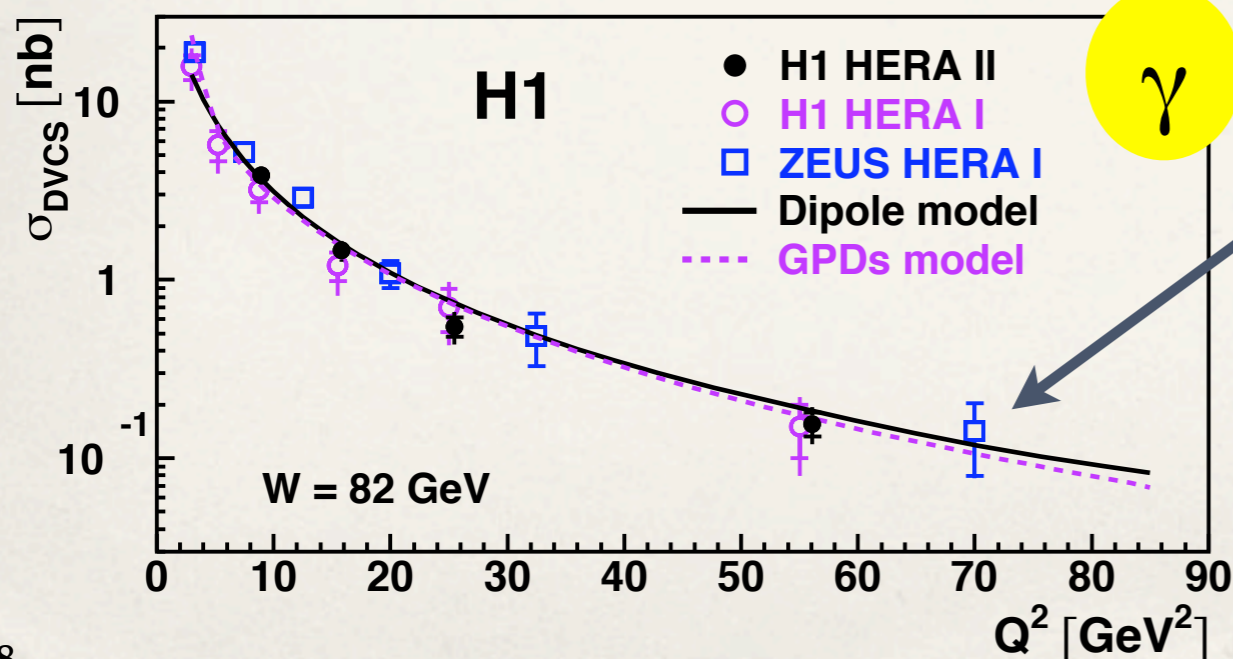
H1 and ZEUS

H1 arXiv:0907.5289



Hard exclusive electroproduction
of a real photon:
Deeply Virtual Compton Scattering,
DVCS

- Steep W^δ dependence
- $\delta \approx 0.7$, independent of Q^2
- DVCS is a hard process: gluons resolved



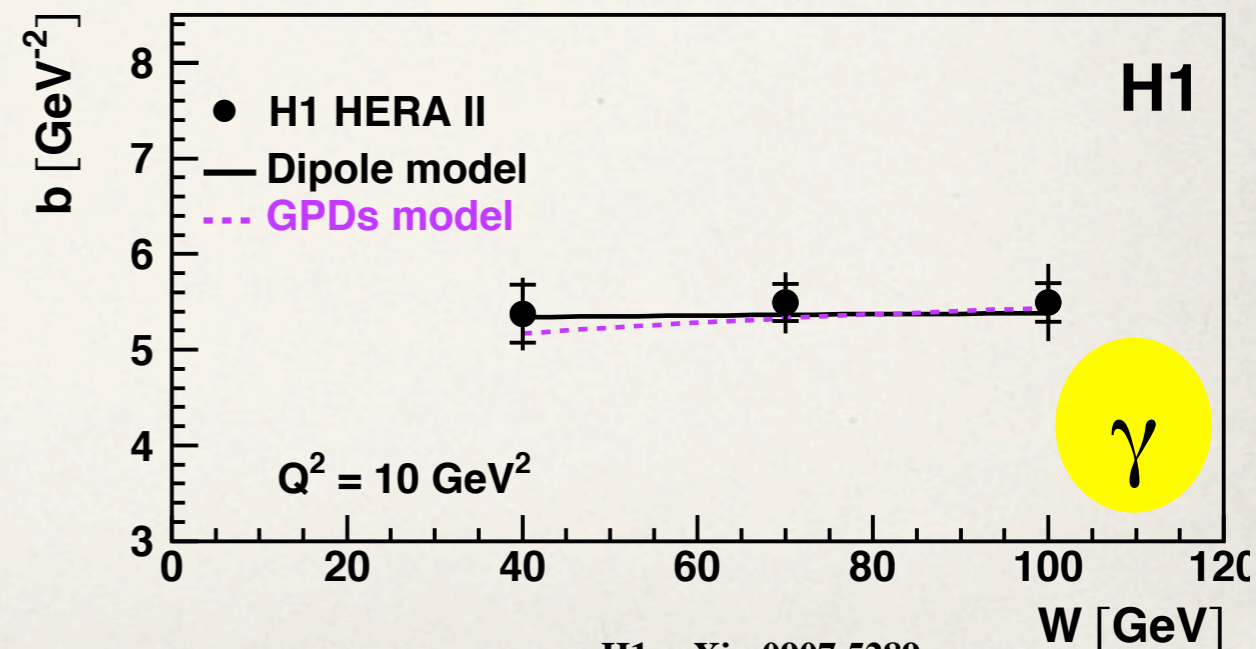
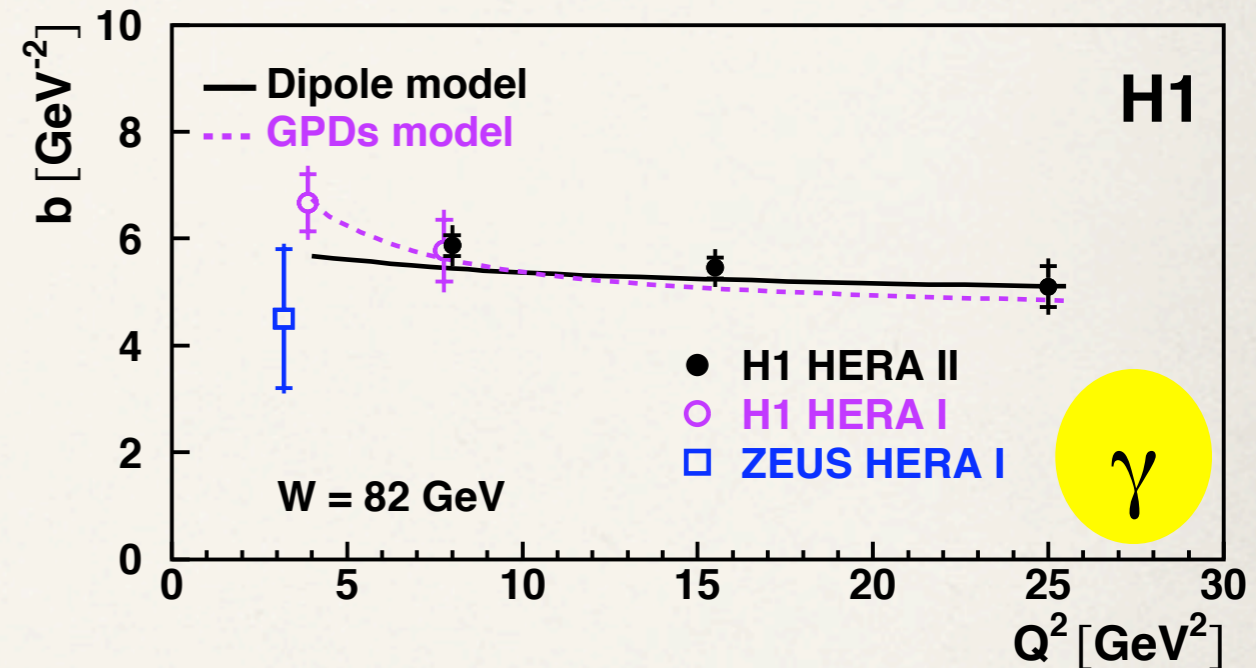
- Decrease of σ_{DVCS} with Q^2
- Q^{2n} dependence, $n \approx 1.5$
- Slower than for vector mesons

Transverse extension of partons

HERA: DVCS cross section differential in t

- Extract $d\sigma/dt$ in bins of Q^2 and W
- Ansatz $d\sigma/dt \propto \exp(-b|t|)$
- t slope \rightarrow average impact parameter b
 - Description of transverse extension of partons
 - Size of proton core (w/o soft periphery)

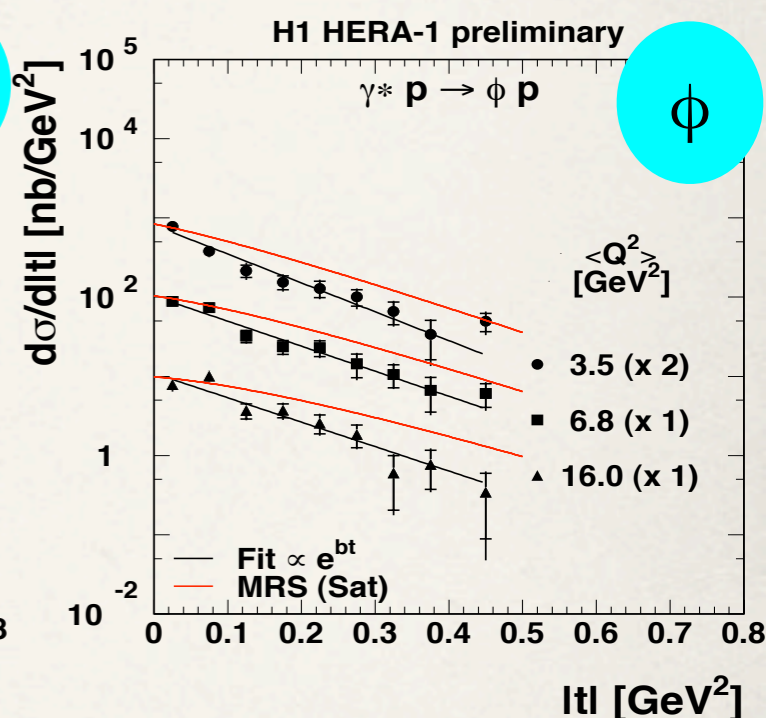
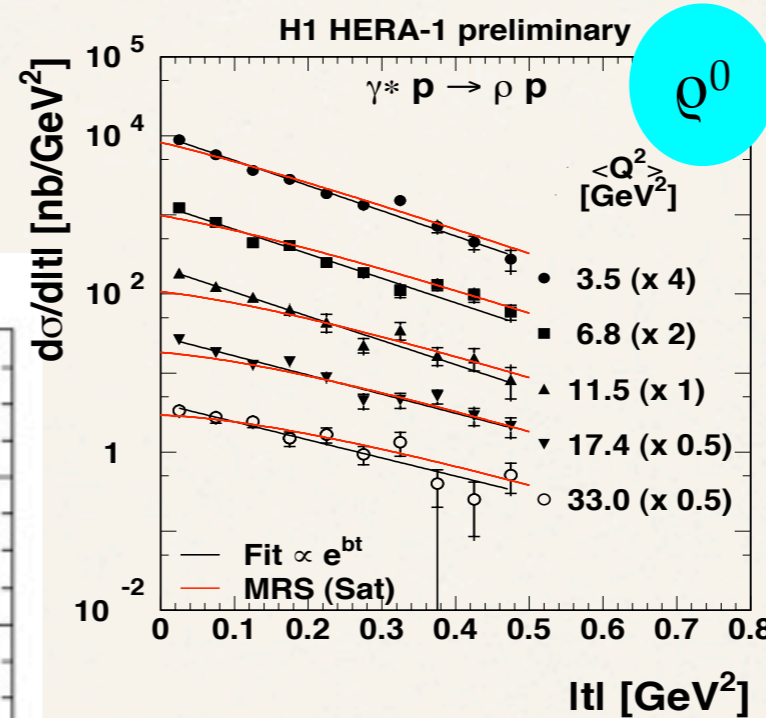
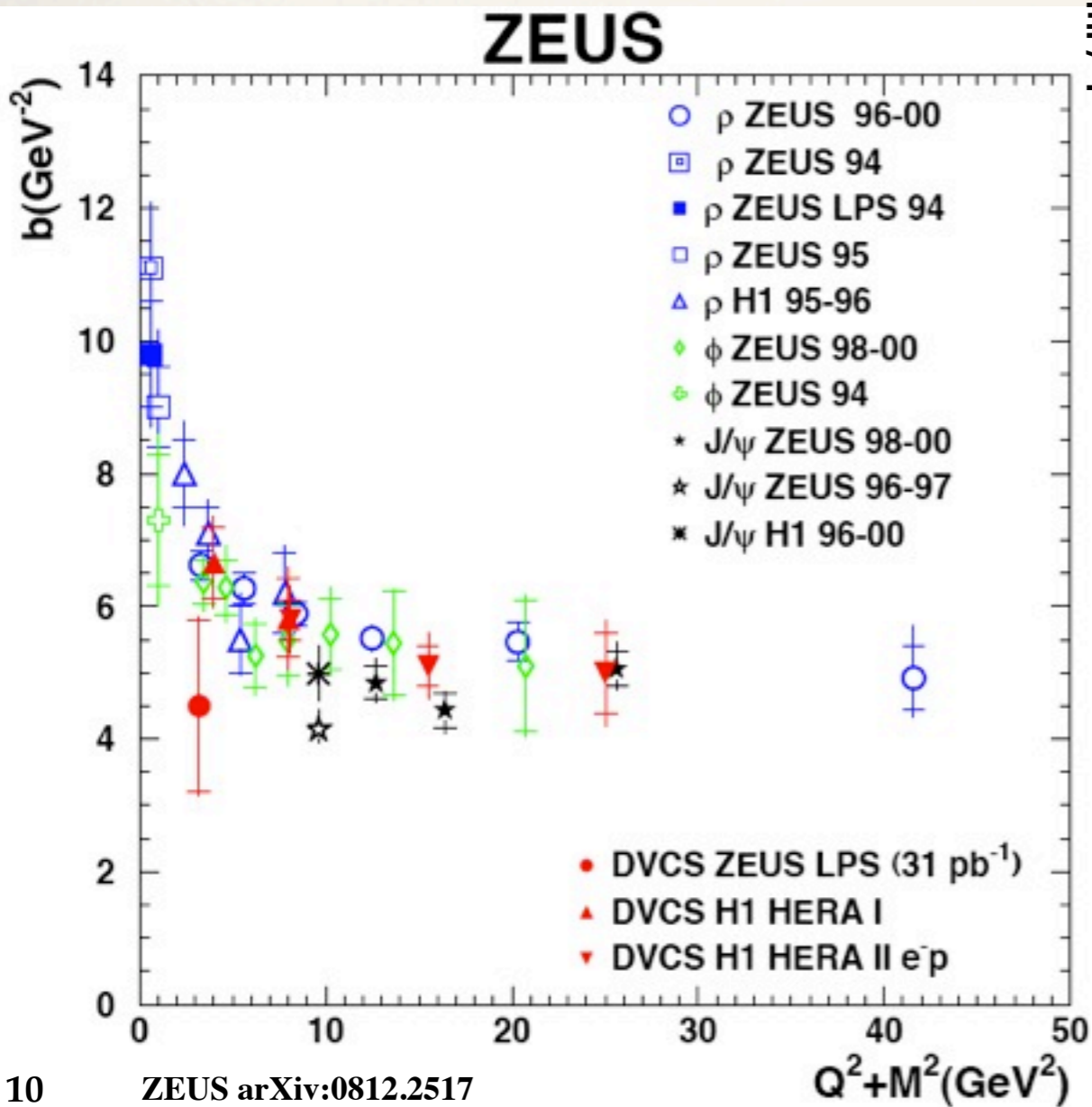
$$\sqrt{\langle r_T^2 \rangle} = (0.65 \pm 0.02) \text{ fm} \quad \text{at } x_B = 10^{-3}$$



H1 arXiv:0907.5289

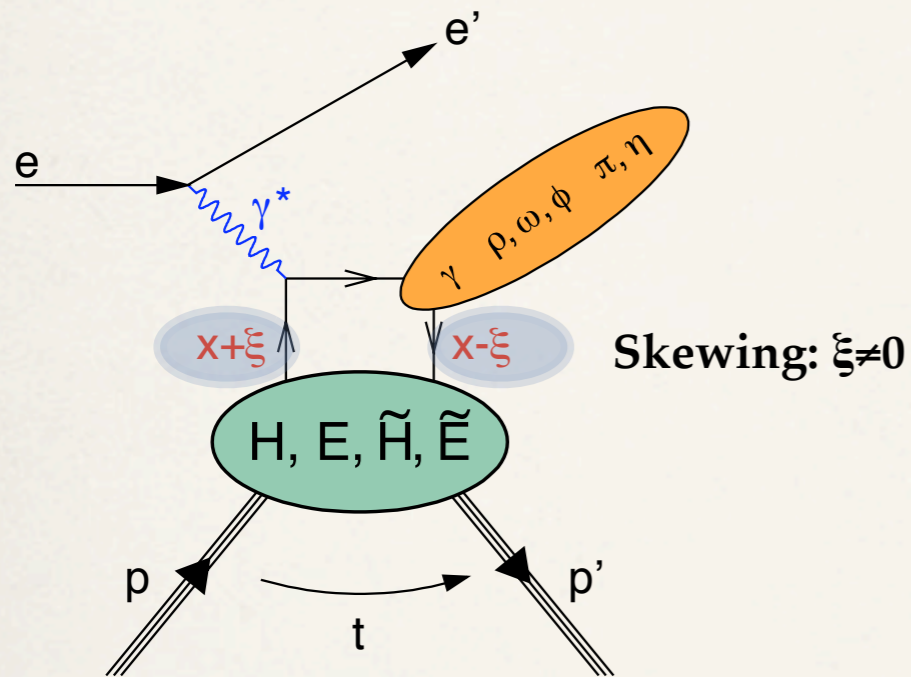
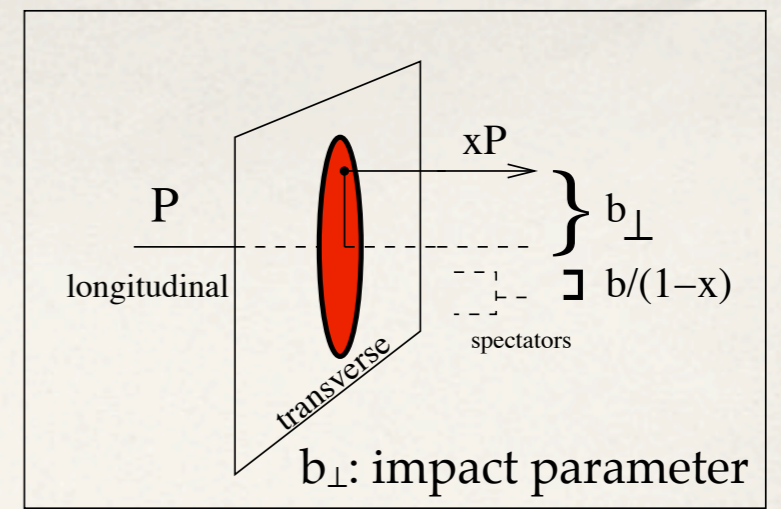
HERA: t slopes

Electroproduction
 $ep \rightarrow epV$



- b measures transverse size of VM \oplus nucleon
- VM shrinks with increasing photon virtuality
- Universal value of $b \approx 5 \text{ GeV}^{-2}$ at large scale

Generalized Parton Distributions



“Nucleon tomography”

PDFs: longitudinal momentum
forward limit $\xi=0, t=0$: $H^q(x, 0, 0) = q(x)$

Form Factors: transverse position
moments of GPDs: $\int_{-1}^1 dx H^q(x, \xi, t) = F_1^q(t)$

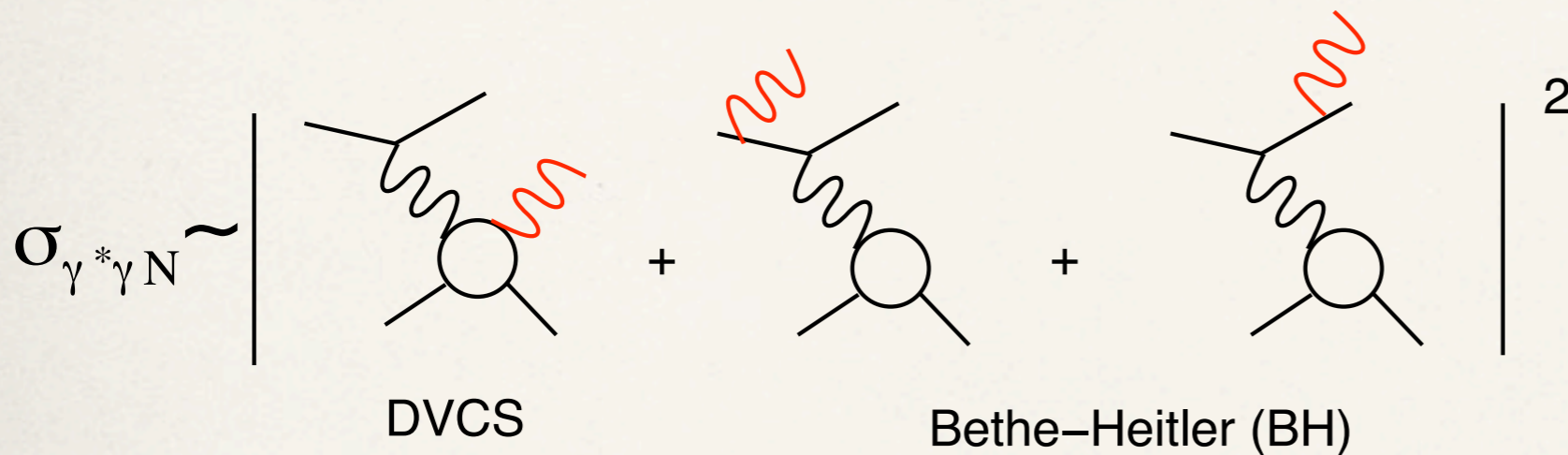
Nucleonic Spin: total angular momentum
Ji relation:

$$J_q = \frac{1}{2} \lim_{t \rightarrow 0} \int_{-1}^1 dx x [H^q(x, \xi, t) + E^q(x, \xi, t)]$$

leading twist, quark chirality conserving, spin-1/2

f(quark helicity)	x	✓
nucleon spin flip	photon: $J^P=1^-$ (DVCS)	
x	H	\tilde{H}
✓	E	\tilde{E}
	$J^P=1^-$ mesons	$J^P=0^-$ mesons

Deeply Virtual Compton Scattering



$$= |\tau_{\text{DVCS}}|^2 + |\tau_{\text{BH}}|^2 + \tau_{\text{DVCS}} \tau_{\text{BH}}^* + \tau_{\text{DVCS}}^* \tau_{\text{BH}}$$

DVCS-BH

interference term \mathcal{I}

Contribution at colliders.

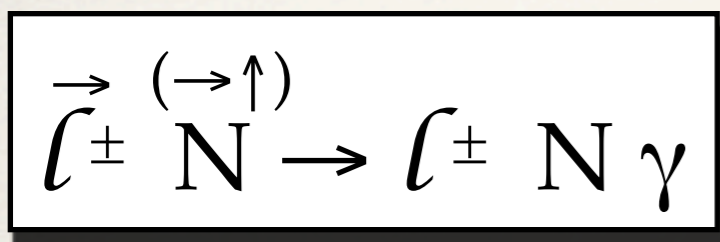
Fixed target:
 $|\tau_{\text{DVCS}}|^2 \ll |\tau_{\text{BH}}|^2$

Exactly calculable in QED
 given the nucleon elastic
 form factors F_1 and F_2

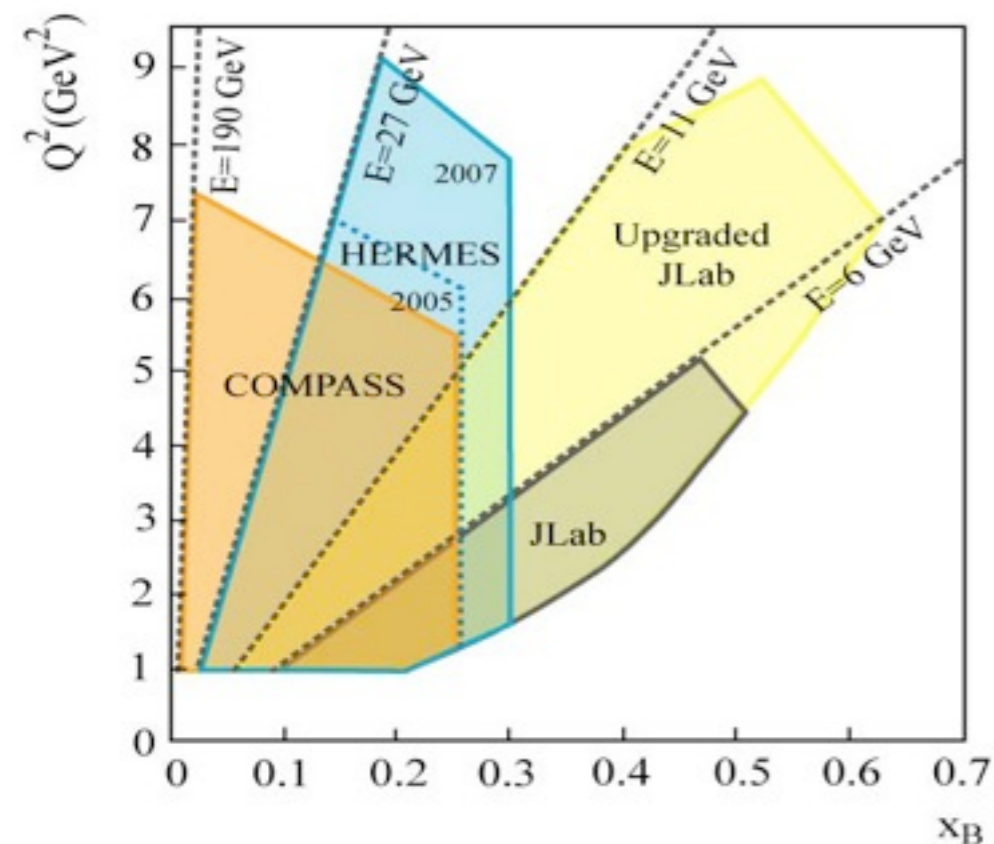
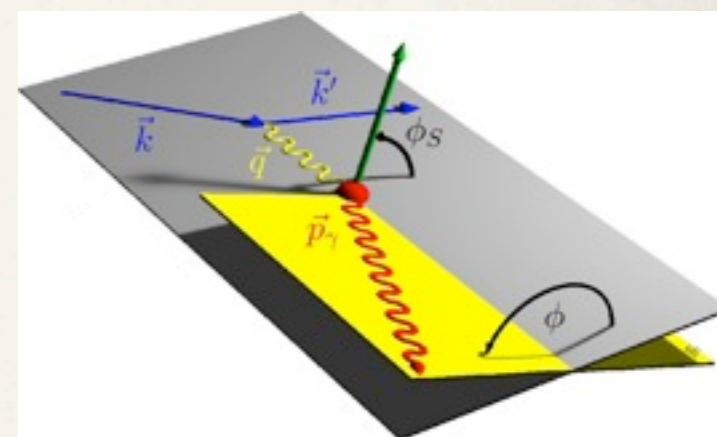
Holographic principle:

- BH reference amplitude magnifies DVCS
- Measure magnitude A and phase φ
 of DVCS amplitude $\tau_{\text{DVCS}} = A e^{i\varphi}$

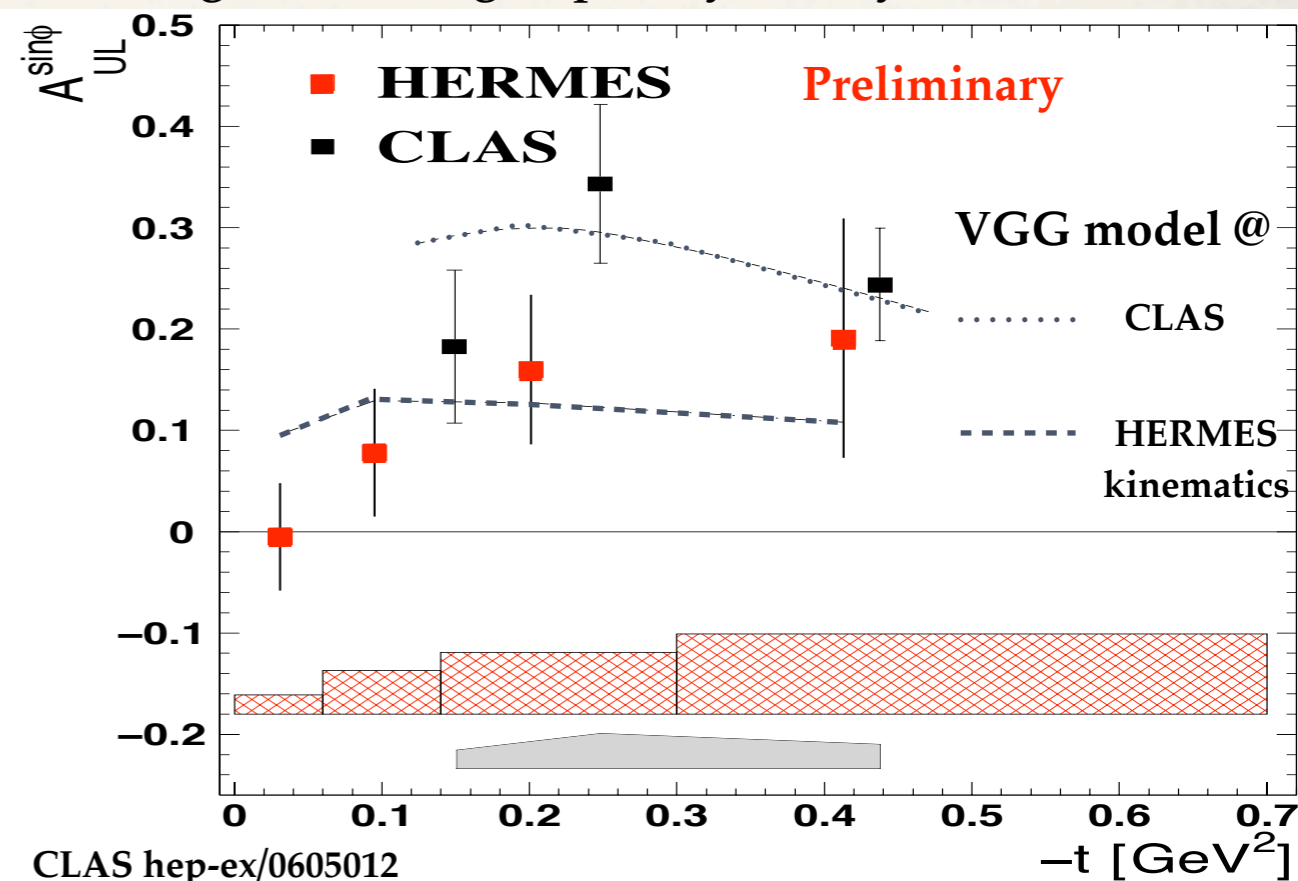
DVCS @ fixed target experiments



The data are subject to a harmonic fit w.r.t. the azimuthal angle(s) ϕ (and ϕ_s) to obtain azimuthal asymmetries

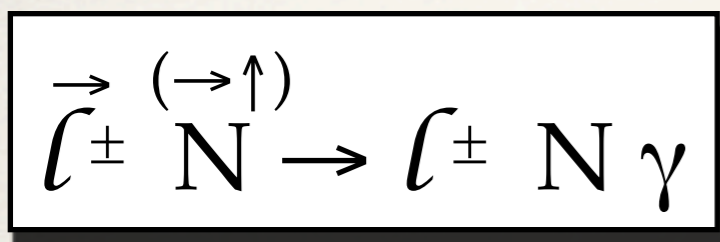


Longitudinal target spin asymmetry in DVCS

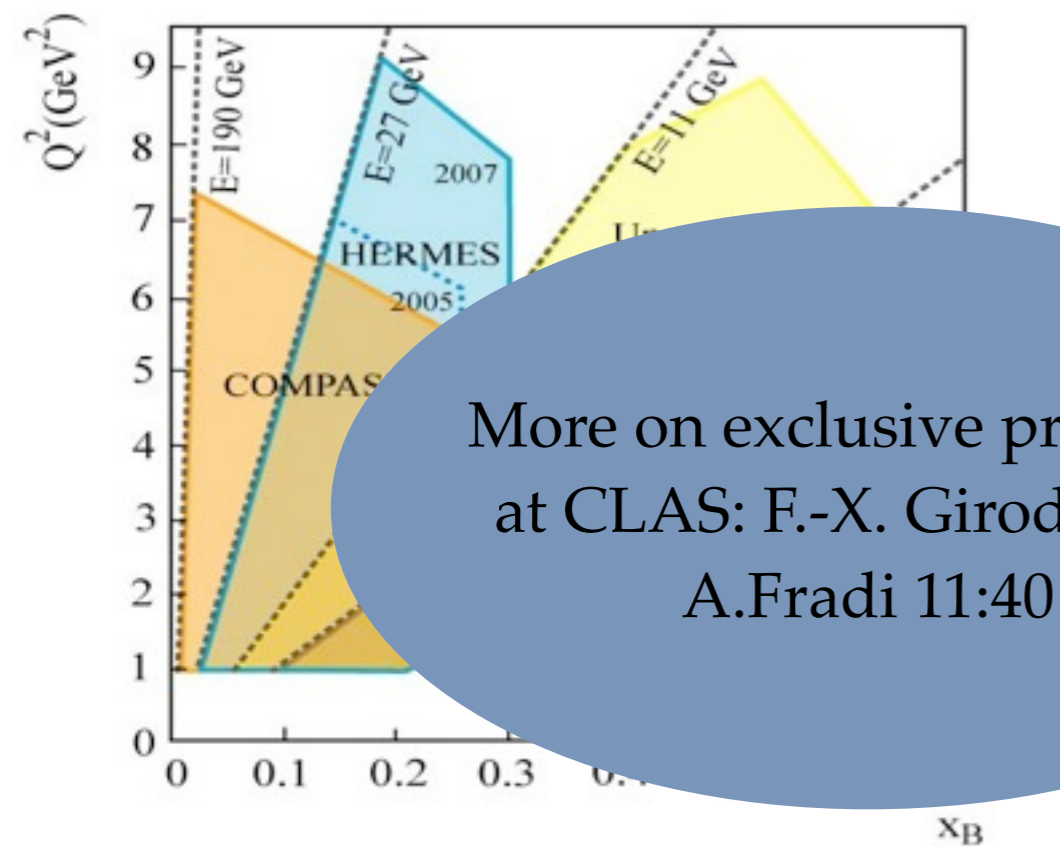
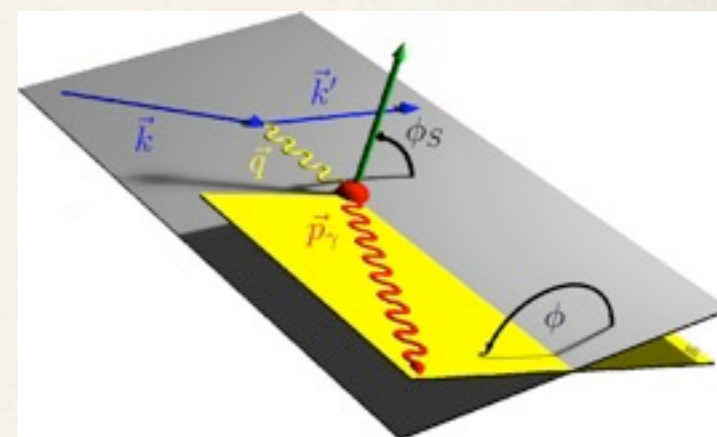


CLAS hep-ex/0605012

DVCS @ fixed target experiments

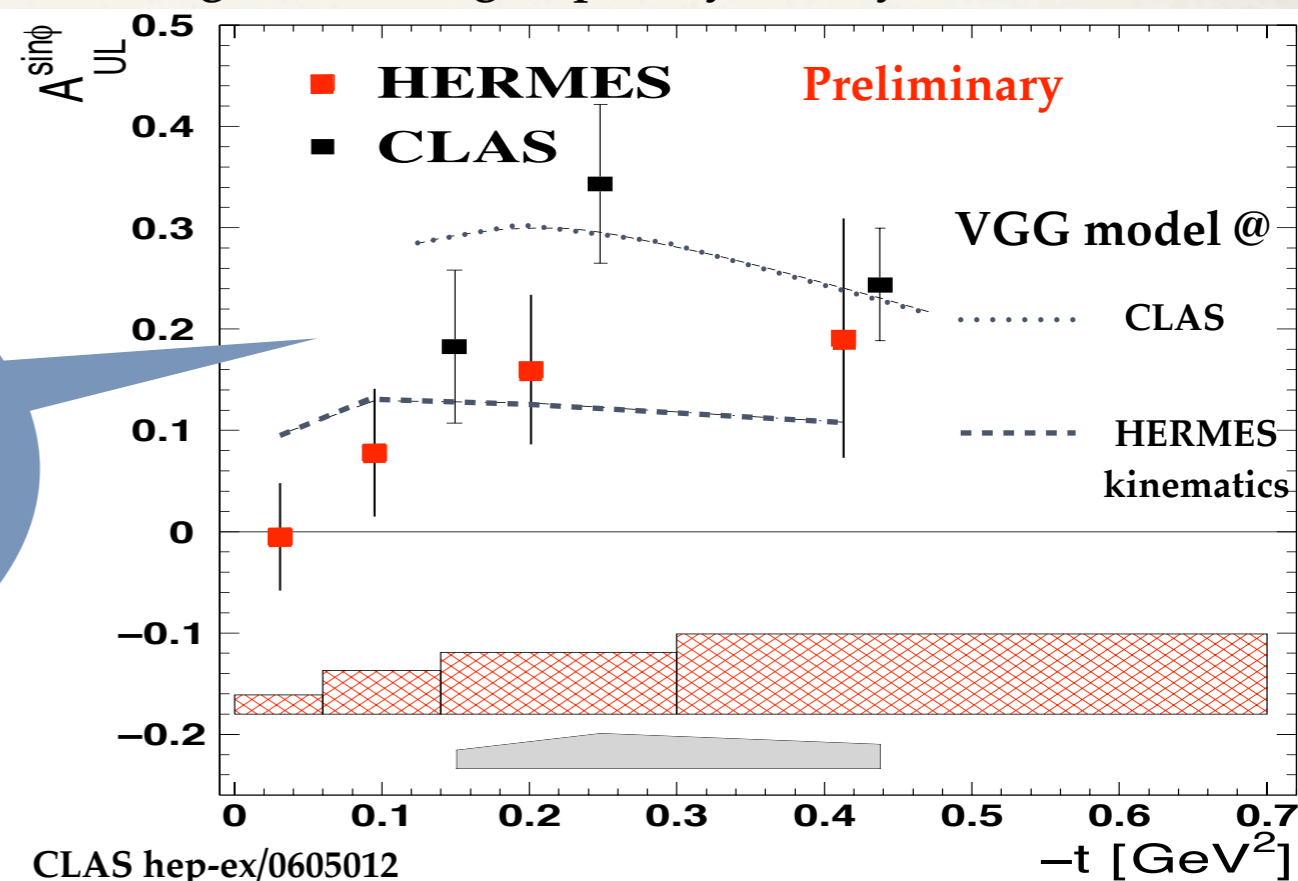


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More on exclusive processes at CLAS: F.-X. Girod 11:00, A.Fradi 11:40

Longitudinal target spin asymmetry in DVCS



CLAS hep-ex/0605012

DVCS cross-section in the valence quark region

- Hall-A at JLab, proton target

Helicity-dependent

$$\propto \text{Im}(\mathcal{T})$$

GPDs @ $x=\xi$

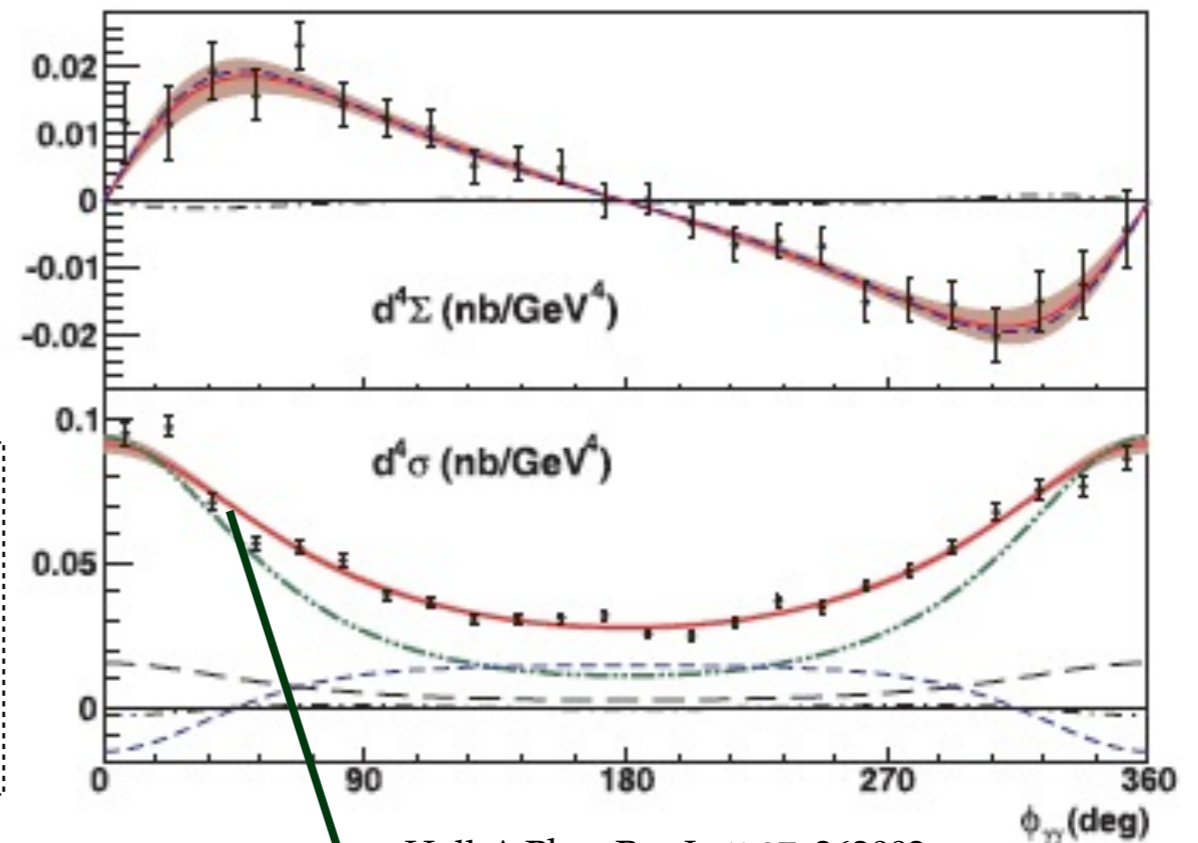
Helicity-independent

$$\propto \text{Re}(\mathcal{T})$$

integral of GPDs over x

Differential cross section vs. azimuthal angle

Bin: $\langle x_B \rangle = 0.36$, $\langle Q^2 \rangle = 2.3 \text{ GeV}^2$, $\langle t \rangle = -0.28 \text{ GeV}^2$



Hall-A Phys.Rev.Lett.97, 262002

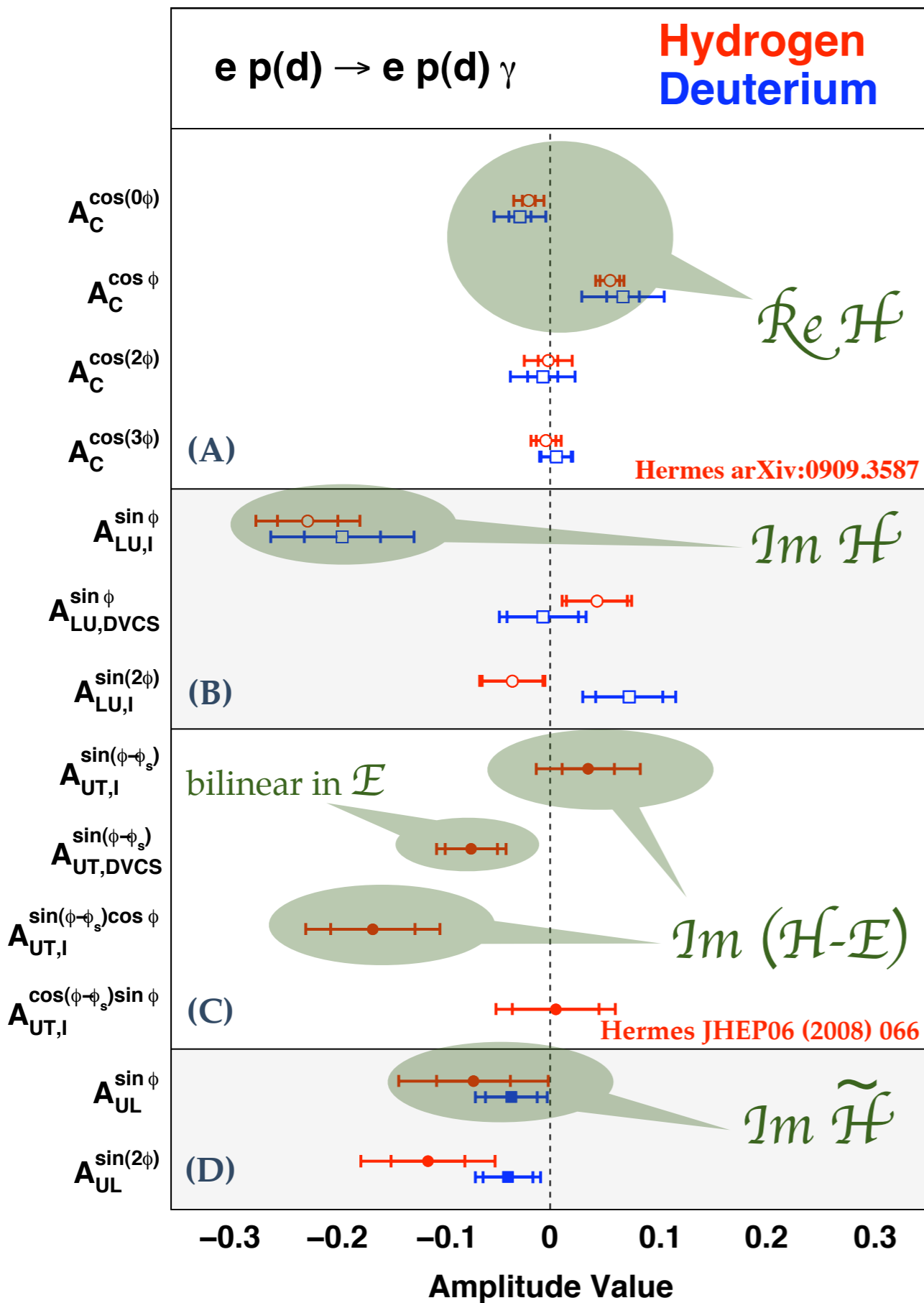
- No Q^2 dependence of $\text{Im}(\mathcal{T})$

- Indication of factorization down to $Q^2 = 2 \text{ GeV}^2$
- GPDs accessible at moderate Q^2

Computed BH contribution
< helicity-independent cross section

DVCS azimuthal amplitudes

HERMES
(prelim.)



(A) Beam charge asymmetry:
GPD H

(B) Beam helicity asymmetry:
GPD H

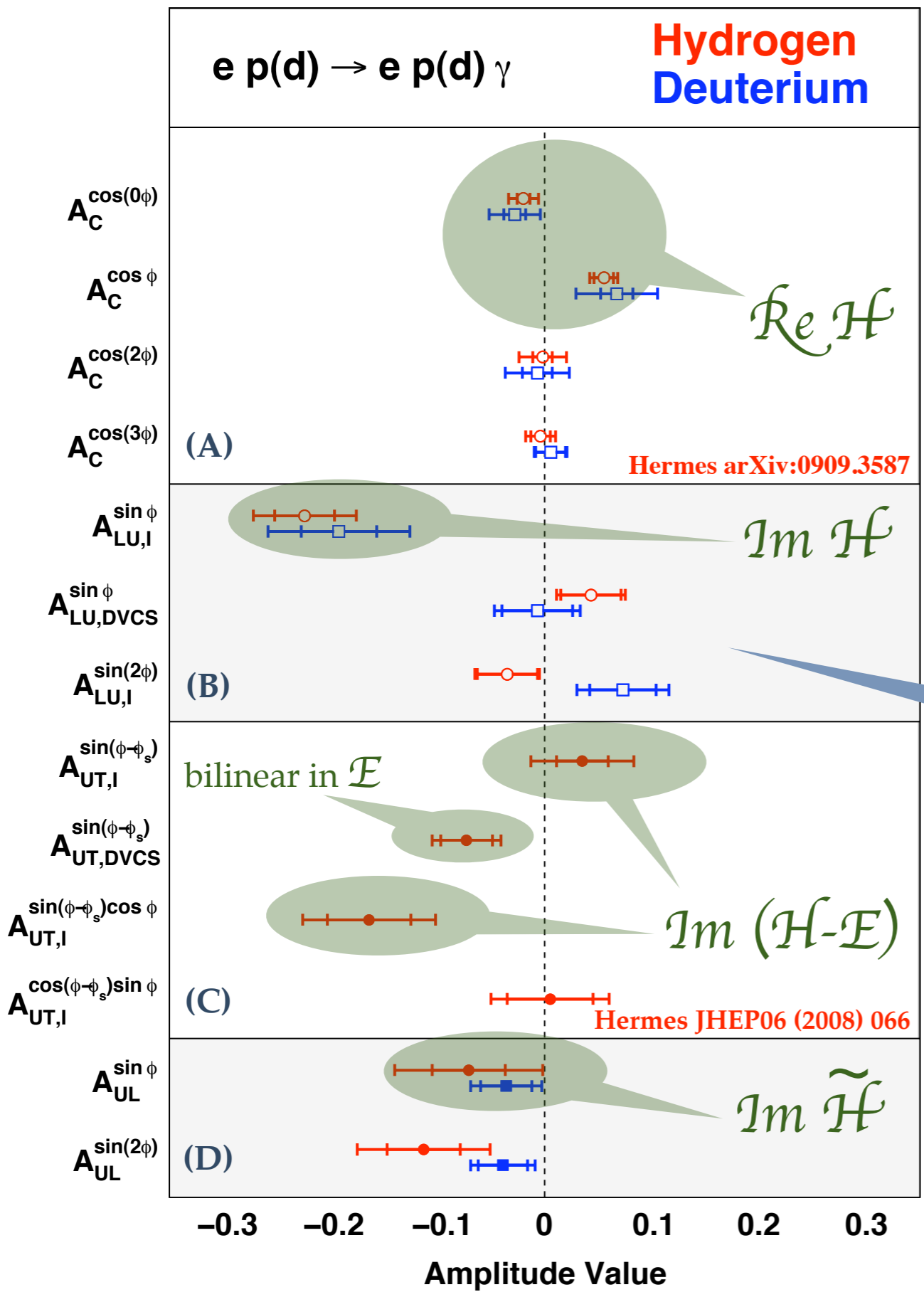
(C) Transverse target spin asymmetry:
GPD E from proton target

(D) Longitudinal target spin asymmetry:
GPD \tilde{H}

$\Re(\tau_{DVCS})$
↑
Projects out
↓
 $\Im(\tau_{DVCS})$

DVCS azimuthal amplitudes

HERMES
(prelim.)



(A) Beam charge asymmetry:
GPD H

(B) Beam helicity asymmetry:
GPD H

(C) Transverse target spin asymmetry:
GPD E and H

(D) Longitudinal target spin asymmetry:
GPD \tilde{H}

$\Re(\tau_{DVCS})$

↑

Projects out

↓

$\Im(\tau_{DVCS})$

Details by
A. Mussgiller 11:20
& HERMES Recoil
detector 2006/07

HERA / H1: Beam Charge Asymmetry in DVCS

First measurement at collider

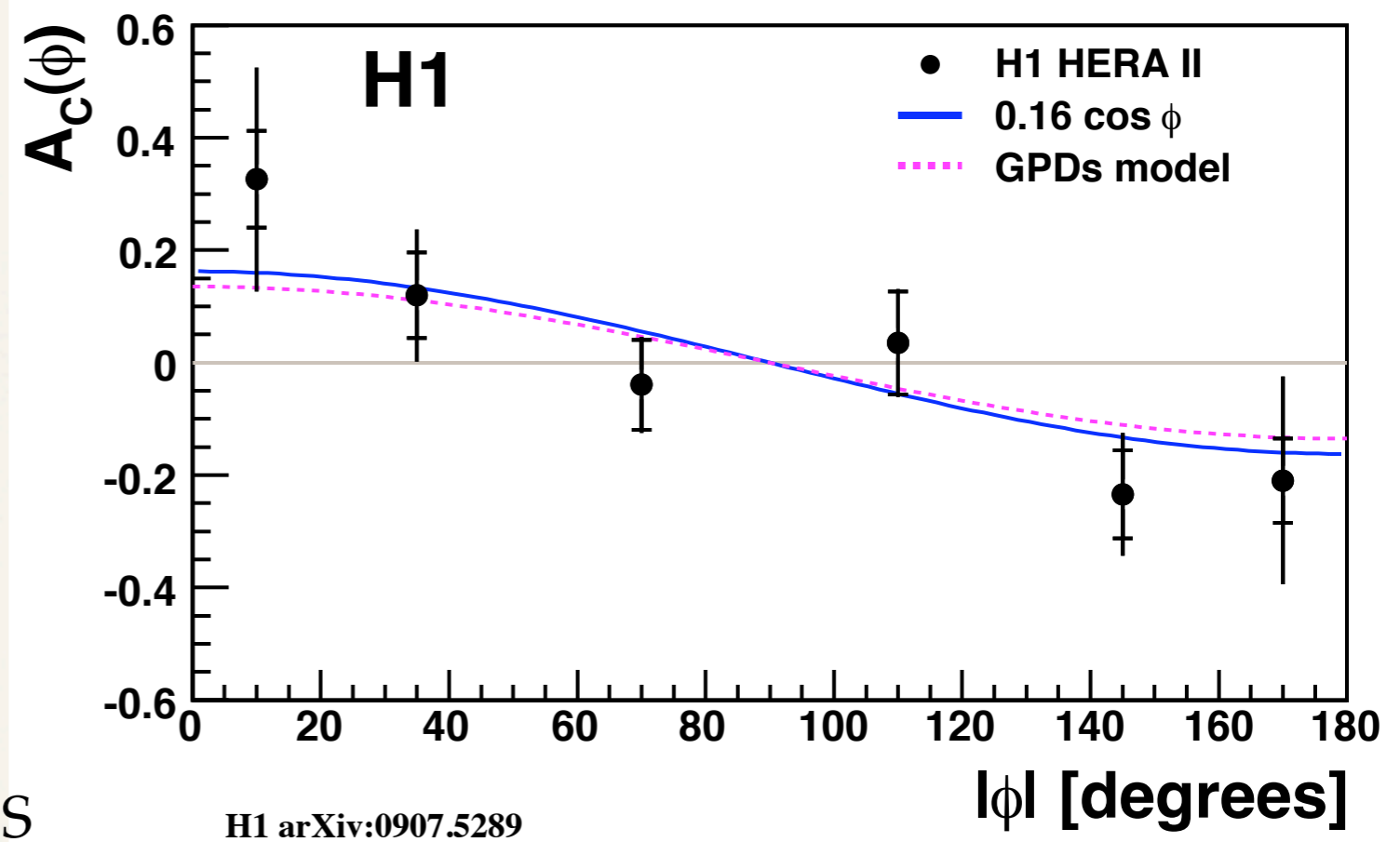
- low $x_B=10^{-4}...10^{-2}$
- $6.5 < Q^2 < 80 \text{ GeV}^2$
- $30 < W < 140 \text{ GeV}$
- $|t| < 1 \text{ GeV}^2$

Positive $\cos\phi$ amplitude

- $\text{Re}(\tau_{\text{DVCS}}) > 0$
- Sign change compared to HERMES

Ratio $\rho = \text{Re}(\tau_{\text{DVCS}}) / \text{Im}(\tau_{\text{DVCS}})$

- $\rho = 0.20 \pm 0.05(\text{stat}) \pm 0.08(\text{sys})$
- In good agreement with calculation from dispersion relation



Access to the total angular momentum of quarks

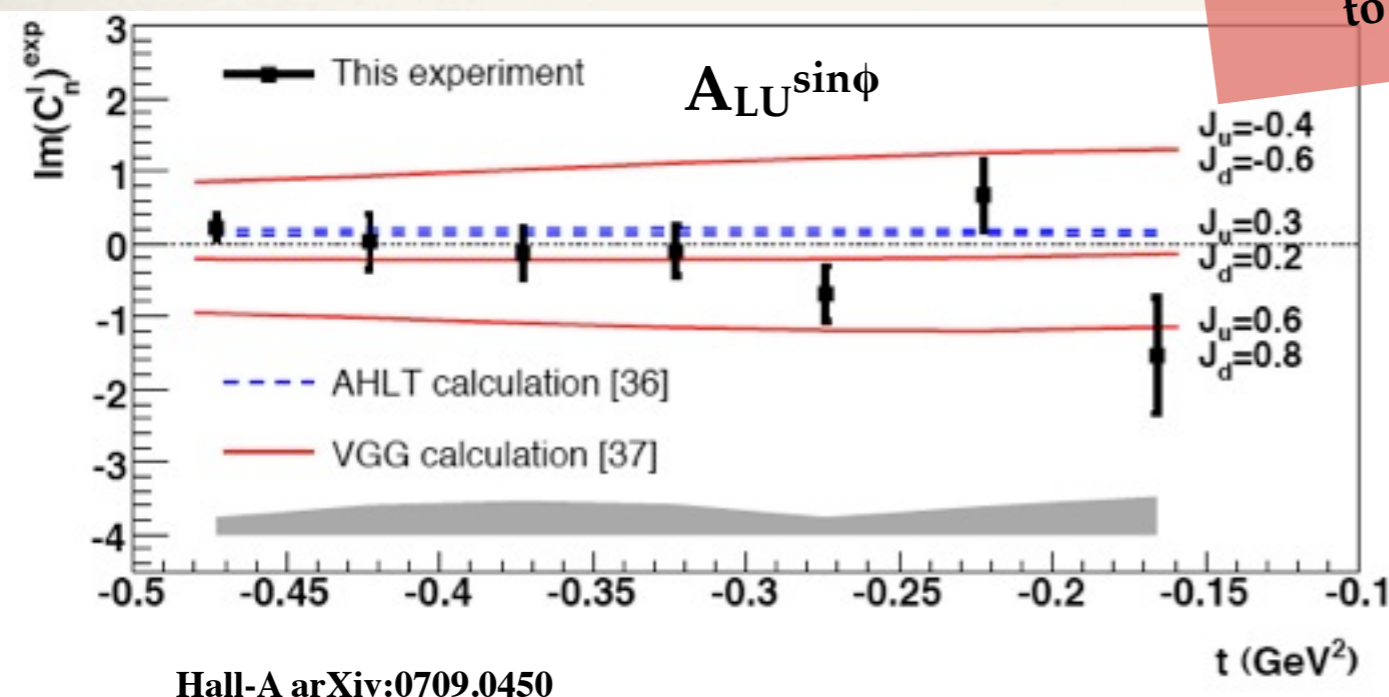
$\vec{e}^- n \rightarrow e^- n \gamma$: sensitive to GPD E_q

Total angular momentum of quarks J_q :

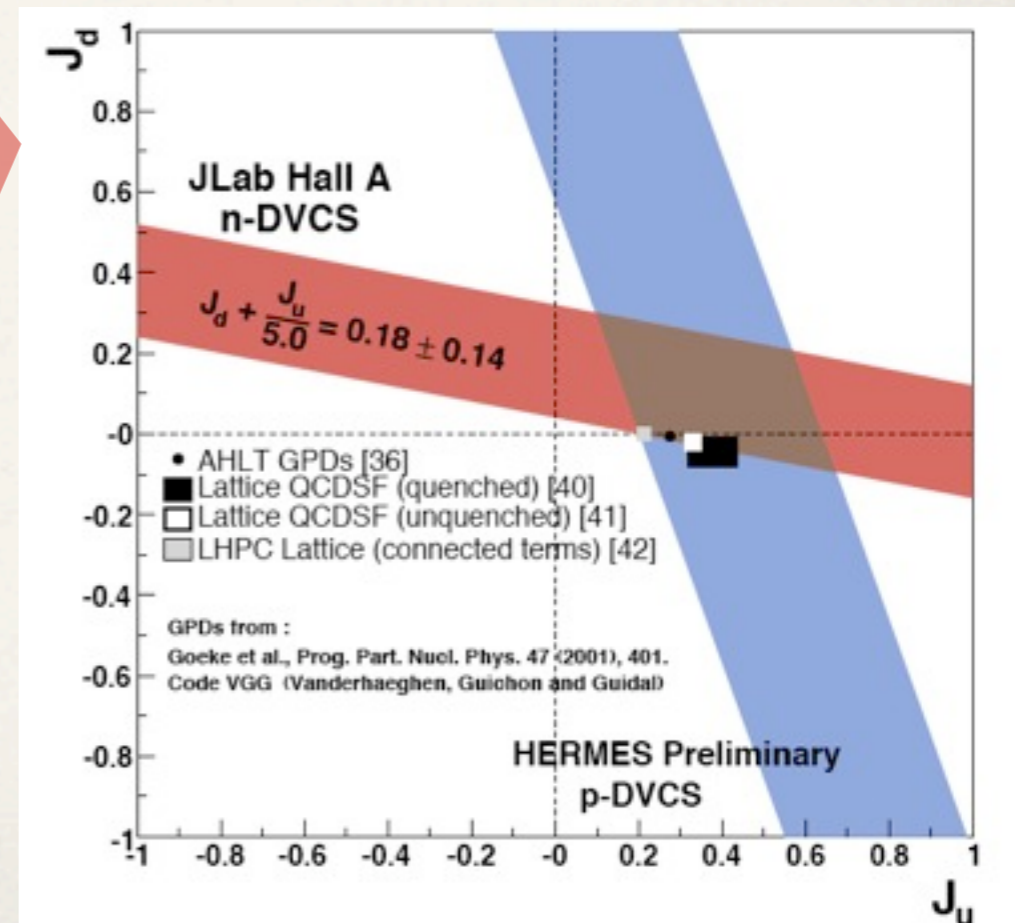
$$J_q = \frac{1}{2} \lim_{t \rightarrow 0} \int_{-1}^1 dx x [H^q(x, \xi, t) + E^q(x, \xi, t)]$$

- Hall-A at JLab, deuteron target
- Quasi-elastic proton contribution subtracted from deuteron signal
- Beam-helicity asymmetry:

from sensitivity to J_q :



Hall-A arXiv:0709.0450



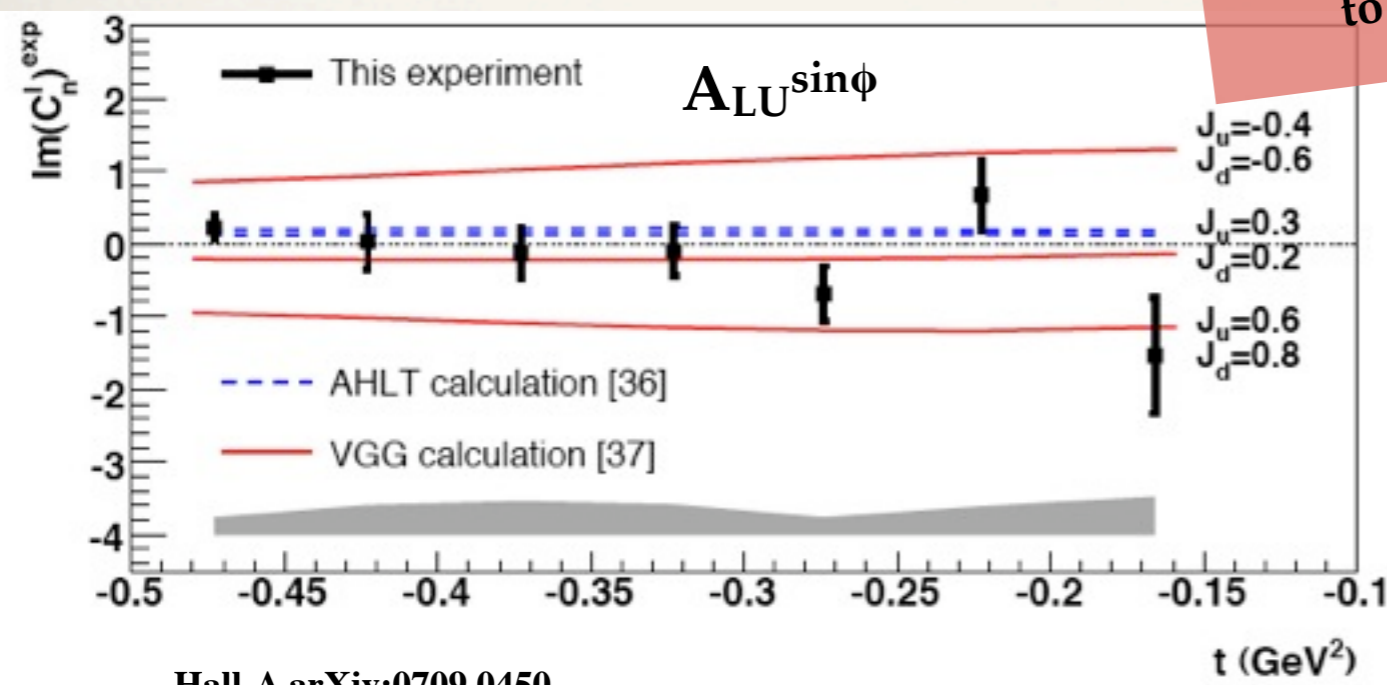
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Total angular momentum of quarks J_q :

$$J_q = \frac{1}{2} \lim_{t \rightarrow 0} \int_{-1}^1 dx x [H^q(x, \xi, t) + E^q(x, \xi, t)]$$

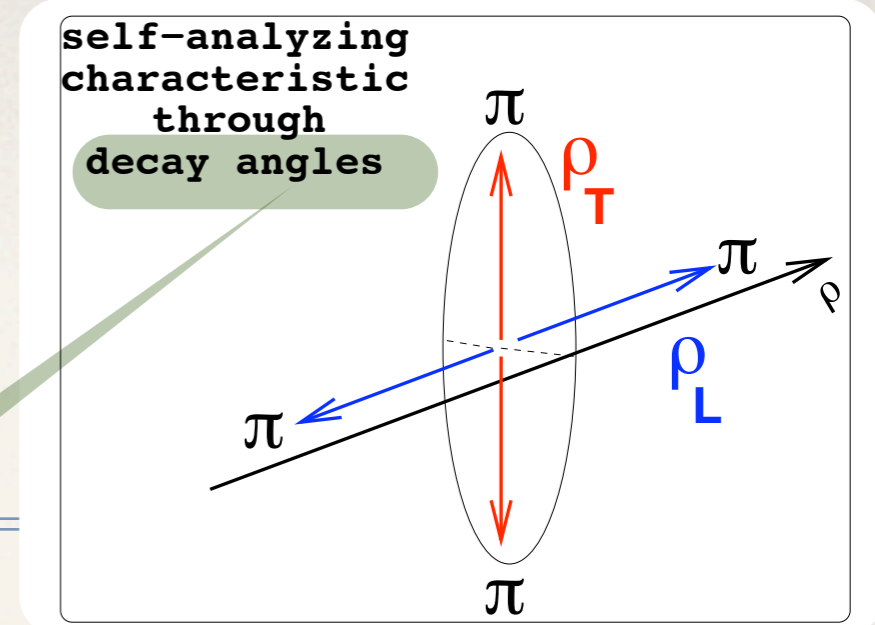


from sensitivity to J_q :

Hall-A arXiv:0709.0450

Caveat: model-dependent constraint on $J_u + k \cdot J_d$.
GPD models are far from describing all available data equally well !!

ρ^0 Spin Density Matrix Elements



- Cross-section for exclusive ρ^0 leptonproduction:

$$\frac{d\sigma}{dx_B dQ^2 dt} W(x_B, Q^2, t, \phi, \phi_S, \varphi, \vartheta)$$

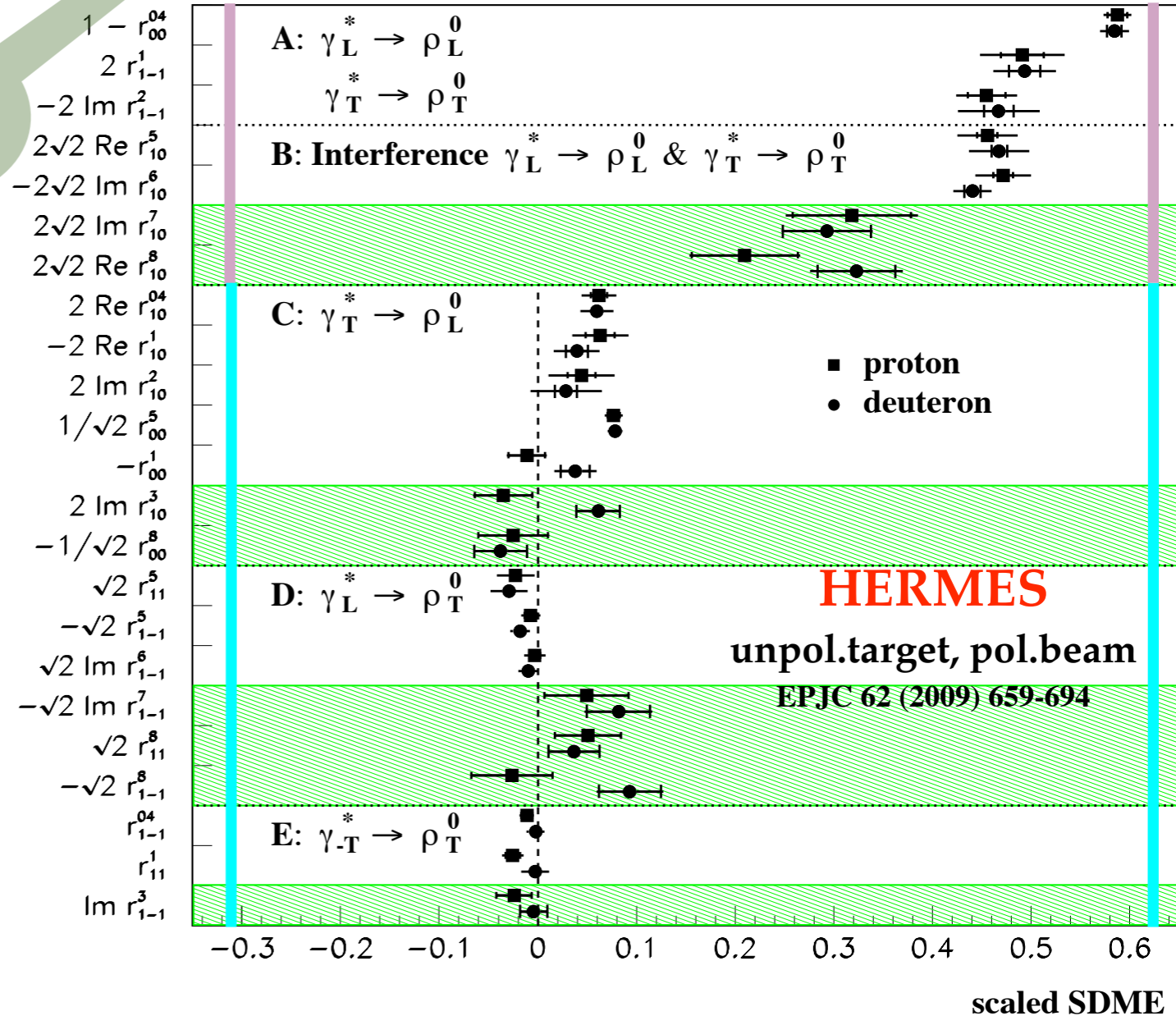
- W parametrized by Spin Density Matrix Elements (SDMEs)

- SDMEs describe helicity transfer from the γ^* to the ρ^0 :

● s-channel helicity conservation (SCHC) $L \rightarrow L, T \rightarrow T$

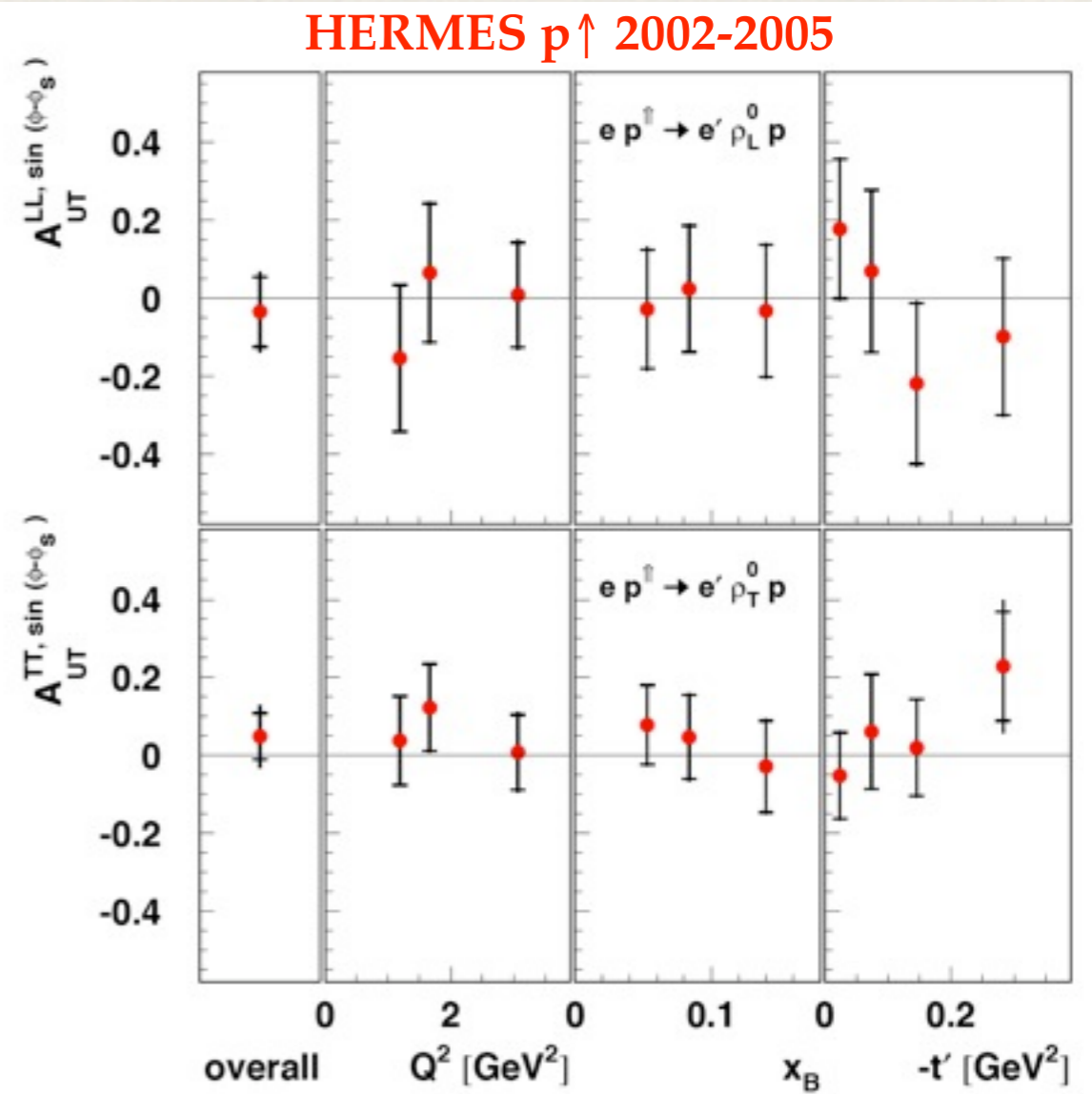
● s-channel helicity violation

- More precision data from HERA

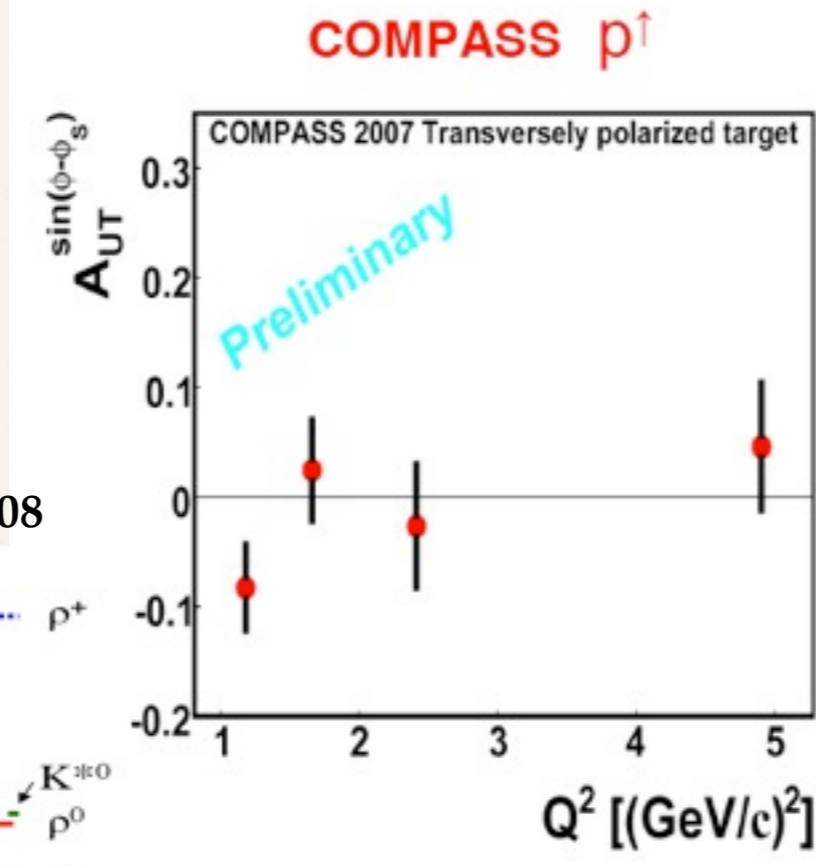


A_{UT} in $ep^{\uparrow} \rightarrow ep\rho^0$

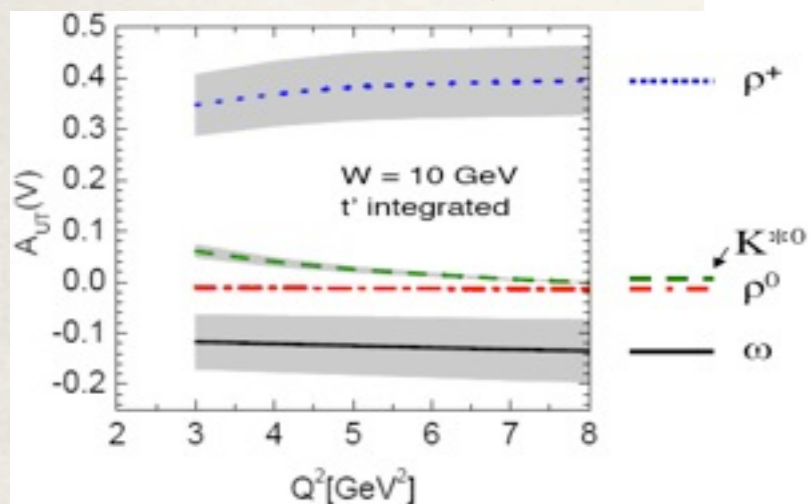
- A_{UT} : depends linearly on GPDs ($E_q + E_g$), no suppression w.r.t. GPD H



HERMES Phys. Lett. B679 (2009) 100-105



GPDs: Goloskokov, Kroll 2008



Exclusive ω A_{UT}

- u- and d-quarks in GPD E do not cancel:
 $A_{UT} \propto \text{Im}[(2E^u - E^d) / (2H^u - H^d)]$
 (≈ -0.10 expected)

- HERMES: $-0.22 \pm 0.16 \pm 0.11$**

The Future

Global GPD fits

See talk by
D. Müller
11:55

COMPASS 2012-15:

large Recoil + polarized target
'DVCS test runs' 08/09, small Recoil

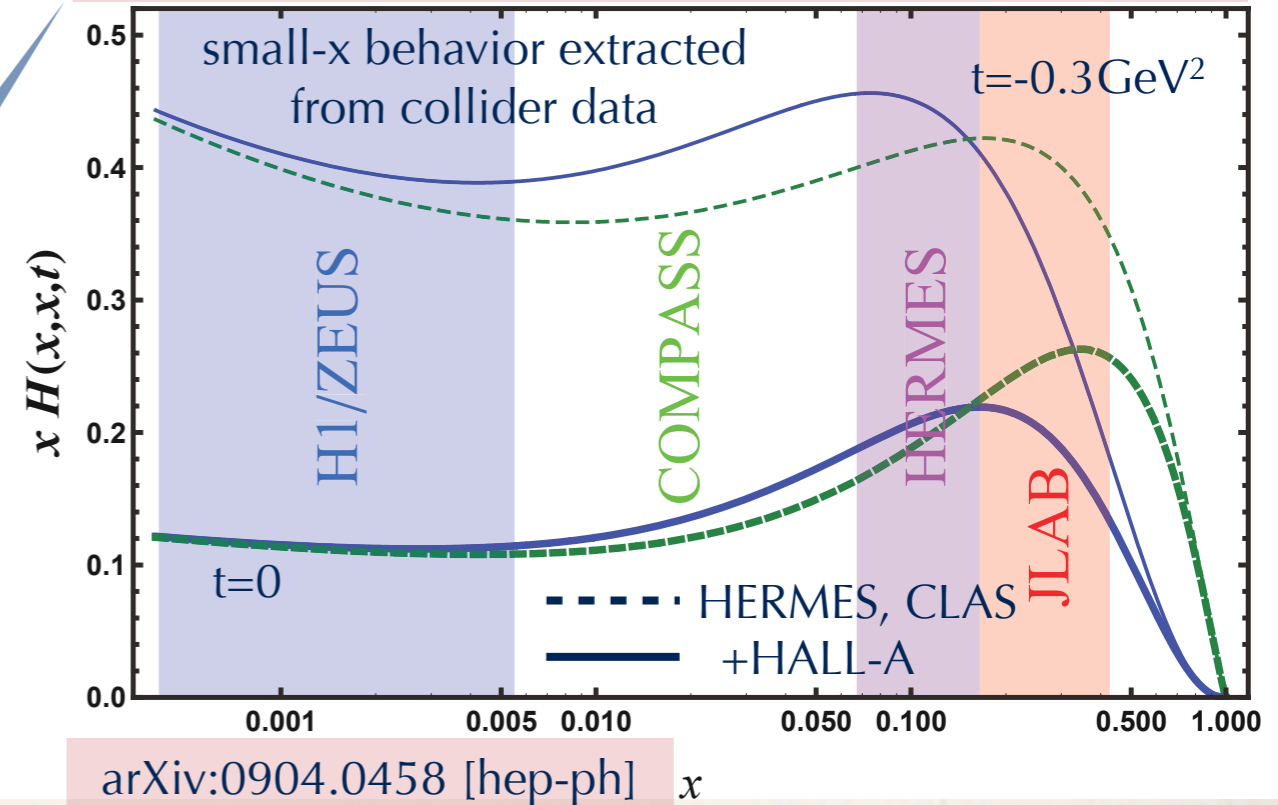
JLab 12 GeV upgrade

$Q^2_{\max} = 13...14 \text{ GeV}^2$, e^+ beam

Future Electron-Ion Collider ($e\vec{p}$ and eA)

eRHIC @ BNL: $\sqrt{s} = 15-200 \text{ GeV}$ (HERMES: 7 GeV),
ENC @ GSI: $\sqrt{s} = 40 \text{ GeV}$, ...

Global fit to $H(x, \xi=x, t)$ from DVCS data



See talk by N. d'Hose
Friday afternoon

See EIC workshop
overview Friday afternoon

Summary:

Hard Exclusive Reactions

- Wide spectrum of measurements of cross sections and azimuthal asymmetries
- Hard exclusive reactions as testing ground for QCD
 - Universal t slope
- Access to Generalized Parton Distributions
 - Different quantum numbers of final state select different GPDs
 - GPDs provide 3-dimensional picture of nucleons
 - GPDs allow (in principle) to constrain total angular momentum of quarks

Thanks to everybody who delivered input to this talk!!